Bismut-Cheeger Eta Form and Higher Spectral Flow

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We prove the variation formula, the embedding formula, and the adiabatic limit formula for equivariant Bismut–Cheeger eta forms with a fiberwise compact Lie group action. We only need the condition that the kernels of corresponding fiberwise Dirac operators form vector bundles.

1 Introduction

The Atiyah–Patodi–Singer eta invariant is the boundary contribution of the index theorem for manifold with boundary with the global boundary condition [2]. Its family extension, the Bismut–Cheeger eta form, is well defined for a fibration of spin manifolds when the kernel of the fiberwise Dirac operator forms a vector bundle [8]. Moreover, the Bismut–Cheeger eta form can be naturally extended to the equivariant case for a fiberwise compact Lie group action. If the base manifold of the fibration is a point and the group is trivial, the equivariant eta form degenerates to the eta invariant. So the equivariant eta form can be considered as the equivariant higher-degree version of the eta invariant.

The purpose of this paper is to generalize some properties of eta invariants: the variation formula, the embedding formula, and the adiabatic limit formula, to

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equivariant Bismut–Cheeger eta forms. For the variation formula and the adiabatic limit formula, we remove additional technical assumptions in previous paper [21].

In [22, 23], the author proved these properties for the equivariant eta form with perturbation, which is the equivariant version of the modified eta form in [30, 31]. In this paper, we obtain these formulas for the equivariant version of the original Bismut–Cheeger eta form by comparing these two versions of equivariant eta forms and the analysis of the equivariant Dai–Zhang higher spectral flow.

1.1 Properties of reduced eta invariants

Let (X, g^{TX}) be an odd dimensional closed spin manifold. Let $\underline{E} := (E, h^E, \nabla^E)$ be a triple that consists of a complex vector bundle E over X, a Hermitian metric h^E on E and a connection ∇^E preserving h^E . We call \underline{E} a geometric triple. Let D_X^E be the Dirac operator twisted with these geometric data. The Atiyah–Patodi–Singer eta invariant of the Dirac operator is defined by

$$\eta\left(D_X^E\right) := \int_0^\infty \operatorname{Tr}\left[D_X^E \exp\left(-u(D_X^E)^2\right)\right] \frac{\mathrm{d}u}{\sqrt{\pi \, u}},\tag{1.1}$$

which is a global spectral invariant. The reduced version

$$\bar{\eta}\left(D_X^E\right) := \frac{1}{2}\eta\left(D_X^E\right) + \frac{1}{2}\dim \ker D_X^E \tag{1.2}$$

appears as the boundary term in the index theorem of Atiyah–Patodi–Singer [2] for compact manifold with boundary.

Note that the eta invariant depends on the geometric data (g^{TX}, h^E, ∇^E) . So it is not a topological invariant. But when the geometric data varies, the variation formula could be written explicitly.

For i = 0, 1, let g_i^{TX} be Riemannian metrics on TX and ∇_i^{TX} be Levi-Civita connections with respect to g_i^{TX} . Let $\underline{E_i} = (E, h_i^E, \nabla_i^E)$, i = 0, 1, be two geometric triples over X. Let D_i be the Dirac operators associated with the geometric data (g_i^{TX}, ∇_i^E) , respectively.

Theorem 1.1. [3] The difference of two reduced eta invariants

$$\bar{\eta} (D_1) - \bar{\eta} (D_0) = \int_X \widetilde{\widehat{A}} \left(TX, \nabla_0^{TX}, \nabla_1^{TX} \right) \operatorname{ch} \left(E, \nabla_0^E \right) + \int_X \widehat{A} \left(TX, \nabla_1^{TX} \right) \widetilde{\operatorname{ch}} \left(E, \nabla_0^E, \nabla_1^E \right) + \operatorname{sf} \left(D_0, D_1 \right). \quad (1.3)$$

Here, $\widehat{A}(\cdot)$, $ch(\cdot)$ are the \widehat{A} -form and the Chern character form [4, Section 1.5], $\widehat{\widehat{A}}(\cdot)$, $\widetilde{ch}(\cdot)$ are corresponding Chern–Simons forms [29, Definition B.5.3] and $sf(D_0, D_1) \in \mathbb{Z}$ is the spectral flow between D_0 and D_1 [3, p. 94].

Another important property of the eta invariant is the Bismut–Zhang embedding formula [11], which plays the key role in the differential index theorem [19] and the equivariant version of which [22] is an important component in the proof of the localization formula of eta invariants [24].

Theorem 1.2. [11] Let $i: Y \to X$ be an embedding between two odd-dimensional closed spin manifolds. Let g^{TX} be a Riemannian metric on TX and g^{TY} be the induced metric on TY. Let ∇^{TY} and ∇^{TX} be corresponding Levi–Civita connections. Let $\underline{\mu} = (\mu, h^{\mu}, \nabla^{\mu})$ be a geometric triple over Y. Then there exist geometric triples $\underline{\xi_{\pm}} = (\xi_{\pm}, h^{\xi, \pm}, \nabla^{\xi, \pm})$ over X satisfying the metric condition [11,(1.10)] and a spectral flow term $\mathrm{sf}(Y, X) \in \mathbb{Z}$ such that

$$\begin{split} \bar{\eta}(D_X^{\xi_+}) &- \bar{\eta}(D_X^{\xi_-}) = \bar{\eta}(D_Y^{\mu}) + \int_X \widehat{A}(TX, \nabla^{TX})\gamma(Y, X) \\ &+ \int_Y \widetilde{A}(TY, \nabla^{TY, N}, \nabla^{TX|_Y}) \widehat{A}^{-1}(N, \nabla^N) \operatorname{ch}(E, \nabla^E) + \operatorname{sf}(Y, X). \end{split}$$
(1.4)

Here, ∇^N is the connection on the normal bundle *N* of *Y* in *X* induced from the connection ∇^{TX} , $\nabla^{TY,N} := \nabla^{TY} \oplus \nabla^N$, and $\gamma(Y, X)$ is the Bismut–Zhang current introduced in [11].

Remark that the vector bundles ξ_{\pm} over X are the Atiyah-Hirzebruch direct image [1] of μ over Y and the geometric triples $\underline{\xi_{\pm}}$ were constructed in [11, (1.10)]. Bismut and Zhang [11] established this embedding formula as an equation mod \mathbb{Z} . In [18, Theorem 4.1], when X and Y are boundaries, the hidden mod \mathbb{Z} term was explained as some APS index. In [22], the author explained the mod \mathbb{Z} term as a spectral flow constructively in general case.

Now, we introduce the adiabatic limit formula of eta invariants established by Bismut and Cheeger [8] and Dai [16]. Let $\pi : W \to B$ be a submersion of two closed manifolds with fiber X. Let $TX := \ker(\pi_* : TW \to TB)$ be the relative tangent bundle. Let T^HW be a horizontal subbundle of TW such that $TW = T^HW \oplus TX$. Let g^{TX} be a metric on TX. Let ∇^{TX} be the connection on TX associated with (T^HW, g^{TX}) defined in (2.2). Assume that TX has a spin structure. Let \underline{E} be a geometric triple over W. Let D_X^E be the fiberwise Dirac operator (see (2.4)). Assume that $\ker D_X^E$ forms a vector bundle over B. Under this assumption, the Bismut–Cheeger eta form $\tilde{\eta}^{BC}(\pi, E) \in \Omega^*(B)$ (non-equivariant version of Definition 2.1) is well defined. Moreover, by [8, (0.6)] and [16, (0.2)],

$$d\tilde{\eta}^{BC}(\pi, E) = \begin{cases} \int_{X} \widehat{A}(TX, \nabla^{TX}) \mathrm{ch}(E, \nabla^{E}) - \mathrm{ch}(\ker D_{X}^{E}, \nabla^{\ker}), & \text{if } \dim X \text{ is even;} \\ \\ \int_{X} \widehat{A}(TX, \nabla^{TX}) \mathrm{ch}(E, \nabla^{E}), & \text{if } \dim X \text{ is } \mathrm{odd.} \end{cases}$$
(1.5)

Here, ∇^{ker} is the connection on the vector bundle $\ker D_X^E$ defined in (2.6). Let g^{TB} be a Riemannian metric on *TB*. Assume that *B* is spin. Then *W* is also spin. For t > 0, let $g_t^{TW} := g^{TX} \oplus t^{-2}\pi^*g^{TB}$, which is a Riemannian metric on *TW*. Let $D_{W,t}^E$ be the Dirac operator associated with (g_t^{TW}, ∇^E) . The following theorem, called the adiabatic limit formula, was established by Bismut and Cheeger [8] when D_X^E is invertible and later extended to the case that $\ker D_X^E$ forms a vector bundle [16]. See also [26, (54)] and [35, Theorem 2.5] for some discussions and applications.

Theorem 1.3. [8, 16] If dim *W* is odd,

$$\lim_{t \to 0} \bar{\eta} \left(D_{W,t}^E \right) \equiv \begin{cases} \int_B \widehat{A}(TB, \nabla^{TB}) \tilde{\eta}^{BC}(\pi, E) + \bar{\eta} \left(D_B^{\ker D_X^E} \right) \mod \mathbb{Z}, & \text{if } \dim X \text{ is even;} \\ \int_B \widehat{A}(TB, \nabla^{TB}) \tilde{\eta}^{BC}(\pi, E) \mod \mathbb{Z}, & \text{if } \dim X \text{ is } \text{odd.} \end{cases}$$
(1.6)

Remark that if *B* is a point and dim *X* is odd,

$$\tilde{\eta}^{BC}(\pi, E) = \frac{1}{2} \eta(D_X^E).$$
(1.7)

Thus, the Bismut-Cheeger eta form can be considered as the higher degree version of the eta invariant.

1.2 Reduced equivariant Bismut-Cheeger eta form

Let *G* be a compact Lie group that acts on *W* and *B* such that $\pi \circ g = \pi$ for any $g \in G$. Then the *G*-action on *B* is trivial. We assume that the *G*-action preserves the splitting $TW = T^H W \oplus TX$ and the spin structure of *TX*. We assume that *E* is *G*-equivariant and g^{TX} , h^E , ∇^E are *G*-invariant. Then the *G*-action commutes with D_X^E .

Assume that ker D_X^E forms a vector bundle over *B*. Under this assumption, the equivariant Bismut–Cheeger eta form $\tilde{\eta}_g^{BC}(\pi, E) \in \Omega^*(B)$ is well defined for any $g \in G$ (Definition 2.1). For the unity $e \in G$, $\tilde{\eta}_e^{BC}(\pi, E) = \tilde{\eta}^{BC}(\pi, E)$. Moreover, by [21, (1.3)] (cf. also

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[16, (0.2)]),

$$d\tilde{\eta}_{g}^{BC}(\pi, E) = \begin{cases} \int_{X^{g}} \widehat{A}_{g}(TX, \nabla^{TX}) \mathrm{ch}_{g}(E, \nabla^{E}) - \mathrm{ch}_{g}(\ker D_{X}^{E}, \nabla^{\ker}), & \text{if } \dim X \text{ is even;} \\ \int_{X^{g}} \widehat{A}_{g}(TX, \nabla^{TX}) \mathrm{ch}_{g}(E, \nabla^{E}), & \text{if } \dim X \text{ is odd.} \end{cases}$$
(1.8)

Here, X^g is the fixed point set of $g \in G$ on X, $\widehat{A}_g(\cdot)$ and $\operatorname{ch}_g(\cdot)$ are the equivariant \widehat{A} -form and the equivariant Chern character form, respectively (see, e.g., [25, Definition 1.3] and [21, (2.44), (2,45)] for the definitions).

As in (1.2), we define the reduced equivariant eta form $\tilde{\tilde{\eta}}_{q}^{BC}(\pi, E)$ by

$$\tilde{\tilde{\eta}}_{g}^{BC}(\pi, E) = \begin{cases} \tilde{\eta}_{g}^{BC}(\pi, E) + \frac{1}{2} \mathrm{ch}_{g}(\ker D_{X}^{E}, \nabla^{\ker}) & \text{if } \dim X \text{ is odd;} \\ \\ \tilde{\eta}_{g}^{BC}(\pi, E), & \text{if } \dim X \text{ is even.} \end{cases}$$
(1.9)

When *B* is a point and dim *X* is odd, as in (1.7), we have

$$\tilde{\tilde{\eta}}_e^{BC}(\pi, E) = \bar{\eta}(D_X^E).$$
(1.10)

So the reduced equivariant Bismut–Cheeger eta form is the equivariant higher-degree version of the reduced eta invariant.

In this paper, we will generalize Theorems 1.1– 1.3 for reduced eta invariants to reduced equivariant Bismut–Cheeger eta forms. For the analogue of these results for holomorphic torsion forms, see [34, Appendix].

For the generalization, the main difficulties are concentrated on the spectral flow terms. In order to generalize these terms, we need to replace the spectral flow terms by the equivariant higher spectral flow terms. Unfortunately, until now, the equivariant higher spectral flow is only well defined when the family index vanishes as an element of the equivariant K-group of the base space (see [17, Definition 1.5] and [23, Definitions 3.7, 3.8]). In general, the assumption that the kernel of the fiberwise Dirac operator forms a vector bundle cannot guarantee that the family index vanishes.

In this paper, we use a trick to overcome this difficulty by adding new vector bundles to make the equivariant family index vanish. In this case, the equivariant eta form with perturbation is well defined (Definition 2.2). Then we establish a comparison formula between it and the reduced equivariant eta form (1.9). And the generalizations follow from the results for equivariant eta forms with perturbation in [23] and the analysis of the higher spectral flow. This trick is inspired by the perturbation techniques in [19, 26].

Although we use the spin condition here, all results in this paper can be naturally extended to the equivariant Clifford module case.

In the following subsections, we will describe our results in more details.

1.3 Variation formula

Let $\pi: W \to B$ be an equivariant submersion of two closed *G*-manifolds with closed fiber *X* such that *G* acts on *B* trivially and the relative tangent bundle *TX* has an equivariant spin structure. Let $T_0^H W$ and $T_1^H W$ be two equivariant horizontal subbundles. Let g_0^{TX} and g_1^{TX} be two *G*-invariant Riemannian metrics on *TX*. Let ∇_0^{TX} and ∇_1^{TX} be connections in (2.2) associated with $(T_0^H W, g_0^{TX})$ and $(T_1^H W, g_1^{TX})$. Let $\underline{E}_i = (E, h_i^E, \nabla_i^E)$, i = 0, 1, be two equivariant versions of geometric triples over *W*. Let $D_{X,0}^E$ and $D_{X,1}^E$ be two fiberwise Dirac operators associated with (g_0^{TX}, ∇_0^E) and (g_1^{TX}, ∇_1^E) , respectively. Assume that $\ker(D_{X,0}^E)$ and $\ker(D_{X,1}^E)$ form vector bundles over *B*. Let $\tilde{\eta}_{g,0}^{BC}(\pi, E)$ and $\tilde{\eta}_{g,1}^{BC}(\pi, E)$ be corresponding reduced equivariant eta forms in (1.9). Let $K_G^*(B)$ be the equivariant topological *K*-group of *B*. Let $ch_g: K_G^*(B) \to H^*(B, \mathbb{C})$ be the equivariant Chern character map (see (2.14) for the definition of ch_q on $K_G^1(B)$) for $g \in G$.

Theorem 1.4. There exists $x \in K_G^*(B)$, such that modulo exact forms on *B*, for $g \in G$,

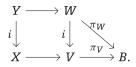
$$\begin{split} \tilde{\eta}_{g,1}^{BC}(\pi, E) &- \tilde{\eta}_{g,0}^{BC}(\pi, E) = \int_{X^g} \widehat{\widehat{A}}_g \left(TX, \nabla_0^{TX}, \nabla_1^{TX} \right) \operatorname{ch}_g \left(E, \nabla_0^E \right) \\ &+ \int_{X^g} \widehat{A}_g \left(TX, \nabla_1^{TX} \right) \widetilde{\operatorname{ch}}_g \left(E, \nabla_0^E, \nabla_1^E \right) + \operatorname{ch}_g(x). \end{split}$$
(1.11)

Here, $\widehat{A}_{g}(\cdot)$, $\widetilde{ch}_{g}(\cdot)$ are equivariant Chern–Simons forms associated with equivariant \widehat{A} -forms and equivariant Chern character forms, respectively, which are the natural equivariant versions of [29, Definition B.5.3].

Remark that if *B* is a point, dim *X* is odd and $G = \{e\}$, Theorem 1.4 degenerates to Theorem 1.1. If $G = \{e\}$, Theorem 1.4 follows directly from the family APS index theorem [30, 31] for the submersion $W \times [0, 1] \rightarrow B$ with fiber $X \times [0, 1]$. In this case, the element $x \in K^*(B)$ in (11) is a family APS index. For the equivariant case, in [21], the author established this formula under the assumption that there exists a smooth path connecting the geometric data $(T_0^H W, g_0^{TX}, h_0^E, \nabla_0^E)$ and $(T_1^H W, g_1^{TX}, h_1^E, \nabla_1^E)$ such that there is no higher spectral flow along this path. Here, we remove this additional assumption. In [23, Theorem 1.2], the author established the variation formula for the equivariant eta forms with perturbation (see Proposition 3.1). For variation formula of eta forms for other settings, see, for example, [12, 13, 20, 26].

1.4 Embedding formula

Let $i : W \to V$ be an equivariant embedding of two closed *G*-manifolds with even codimension. Let $\pi_V : V \to B$ be an equivariant submersion with closed fiber *X*, whose restriction $\pi_W : W \to B$ is also an equivariant submersion with closed fiber *Y*. We assume that *G* acts on *B* trivially. We have the diagram of fibrations:



Let $T^H V$ be an equivariant horizontal subbundle over V. Assume that $T^H V|_W \subset TW$. Set $T^H W := T^H V|_W$. Then $T^H W$ is an equivariant horizontal subbundle over W. Let g^{TX} be an equivariant metric on TX and g^{TY} be the induced metric on TY. We assume that TY and TX have equivariant spin structures. Let N be the normal bundle of TY in TX. Let ∇^{TY} and ∇^{TX} be connections in (2.2) associated with $(T^H W, g^{TY})$ and $(T^H V, g^{TX})$, respectively. Let ∇^N be the connection on N induced from ∇^{TX} . Set $\nabla^{TY,N} := \nabla^{TY} \oplus \nabla^N$.

Let $\underline{\mu} = (\mu, h^{\mu}, \nabla^{\mu})$ be an equivariant geometric triple over *W*. As in Theorem 1.2, by [22, Section 3.3] (see also [11, (1.10)], [24, Section 1.4]), we can construct the equivariant version of the Atiyah–Hirzebruch direct image: the equivariant geometric triples $\underline{\xi_{\pm}} = (\xi_{\pm}, h^{\xi_{\pm}}, \nabla^{\xi_{\pm}})$ over *V* satisfying the equivariant metric condition [22, (3.13)]. Assume that $\ker D_X^{\xi_+}$, $\ker D_X^{\xi_-}$, and $\ker D_Y^{\mu}$ form vector bundles over *B*.

Theorem 1.5. There exists $x \in K_G^*(B)$ such that modulo exact forms on *B*, for $g \in G$,

$$\begin{split} \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{V},\xi_{+}) &- \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{V},\xi_{-}) = \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{W},\mu) + \int_{X^{g}} \widehat{A}_{g}(TX,\nabla^{TX})\gamma_{g}(Y^{g},X^{g}) \\ &+ \int_{Y^{g}} \widetilde{\tilde{A}}_{g}(TY,\nabla^{TY,N},\nabla^{TX|_{W^{g}}}) \widehat{A}_{g}^{-1}(N,\nabla^{N}) \mathrm{ch}_{g}(E,\nabla^{E}) + \mathrm{ch}_{g}(x). \end{split}$$
(1.12)

Here, $\gamma_g(Y^g, X^g)$ is the equivariant Bismut–Zhang current defined in [22].

Remark that for Clifford module case, Theorem 1.5 needs an additional assumption that N has an equivariant spin^c structure as in [22].

When B is a point, dim X is odd and $G = \{e\}$, Theorem 1.5 degenerates to Theorem 1.2. Remark that in [22], the author obtained the equivariant version of Theorem 1.2 and generalized it to the equivariant eta forms with perturbation.

1.5 Functoriality

Let W, V, B be closed G-manifolds. Let $\pi_1 : W \to V$, $\pi_2 : V \to B$ be equivariant submersions with closed fibers X, Y. Then $\pi_3 = \pi_2 \circ \pi_1 : W \to B$ is an equivariant submersion with closed fiber Z. Assume that G acts on B trivially. We have the diagram of fibrations:

X —	$\rightarrow Z -$	$\longrightarrow W$	
	π_1	$\pi_1 \mid \overset{\pi_3}{\checkmark}$	
	1	$\downarrow \pi_2$	k
	Y -	$\longrightarrow V \xrightarrow{2}$	<i>→ B</i> .

Assume that relative tangent bundles TX and TY have equivariant spin structures. So is the relative tangent bundle $TZ \simeq \pi_1^*TY \oplus TX$. Let $(T_1^HW, g^{TX}), (T_2^HV, g^{TY})$ and (T_3^HW, g^{TZ}) be equivariant geometric data with respect to π_1, π_2 , and π_3 as in Section 1.2. Let ∇^{TX}, ∇^{TY} , and ∇^{TZ} be the corresponding connections on TX, TY, and TZ as in (2.2). Set $\nabla^{TY,TX} := \pi_1^* \nabla^{TY} \oplus \nabla^{TX}$. Let $\underline{E} = (E, h^E, \nabla^E)$ be an equivariant geometric triple over W.

Let D_X^E and D_Z^E be fiberwise Dirac operators associated with (g^{TX}, ∇^E) and (g^{TZ}, ∇^E) . Assume that ker D_X^E (resp. ker D_Z^E) forms a vector bundle over V (resp. B). Let ∇^{ker} be the induced G-invariant connection on the vector bundle ker D_X^E as in (2.6). Let $D_Y^{\text{ker}D_X^E}$ be the fiberwise Dirac operator twisted with the vector bundle ker D_X^E over V associated with $(g^{TY}, \nabla^{\text{ker}})$ such that ker $D_Y^{\text{ker}D_X^E}$ forms a vector bundle over B.

Theorem 1.6. There exists $x \in K_G^*(B)$, such that modulo exact forms on B, for $g \in G$, if dim X is even,

$$\begin{split} \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{3}, E) &= \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{2}, \ker D_{X}^{E}) + \int_{Z^{g}} \widetilde{\widehat{A}}_{g}(TZ, \nabla^{TY, TX}, \nabla^{TZ}) \operatorname{ch}_{g}(E, \nabla^{E}) \\ &+ \int_{Y^{g}} \widehat{A}_{g}(TY, \nabla^{TY}) \, \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{1}, E) + \operatorname{ch}_{g}(x); \end{split}$$
(1.13)

if dim X is odd,

$$\begin{split} \tilde{\eta}_{g}^{BC}(\pi_{3}, E) &= \int_{Z^{g}} \widetilde{\widehat{A}}_{g}(TZ, \nabla^{TY, TX}, \nabla^{TZ}) \operatorname{ch}_{g}(E, \nabla^{E}) \\ &+ \int_{Y^{g}} \widehat{A}_{g}(TY, \nabla^{TY}) \, \tilde{\eta}_{g}^{BC}(\pi_{1}, E) + \operatorname{ch}_{g}(x). \end{split}$$
(1.14)

Remark that when *B* is a point, $G = \{e\}$, Theorem 1.6 is a generalization of Theorem 1.3. In fact, let ∇_t^{TZ} be the Levi-Civita connection associated with $g^{TX} \oplus t^{-2}\pi_1^*g^{TY}$. Then, by [21, Proposition 4.5] (cf. also [27, (4.32)]), $\lim_{t\to 0} \widetilde{A}_g(TZ, \nabla^{TY,TX}, \nabla_t^{TZ}) = 0$. The hidden mod \mathbb{Z} term in (1.6) is explained here as a spectral flow.

In [21], we established this formula under the assumptions that $\ker D_Y^{\ker D_X^E} = 0$ and there is no higher spectral flow in all deformations. Here, we remove these assumptions. Notice that if dim X is odd, (14) seems different from [21, Theorem 1.3]. But they are the same under the assumptions in [21, Theorem 1.3]. The reason is that if dim X is odd, the term $\tilde{\eta}_g(T_2^H V, g^{TY}, h^{L_Y}, h^{\ker D^X}, \nabla^{L_Y}, \nabla^{\ker D^X})$ in [21, (1.6)] vanishes by counting degrees on both sides. In [23], we obtain the functoriality for equivariant eta forms with perturbation. See also [5, 13–15, 20, 27, 28] for other settings.

Note that when the fiberwise Dirac operator D is a fiberwise signature operator, ker D naturally forms a vector bundle. In this case, the mod \mathbb{Z} term in (1.6) was constructed by spectral sequences in [16]. When the formulas are extended to the equivariant family case, the spectral sequence terms should be generalized to the equivariant Chern characters of some vector bundles arising from the study of spectral sequence as in [13, 28].

Naturally, the functoriality can be extended to the multifibration case in [8, Appendix 2].

1.6 The organization of the article

This article is organized as follows. In Section 2, we recall the definition of the equivariant Bismut-Cheeger eta form and compare it with the equivariant eta form with perturbation. In Sections 3-5, we prove Theorems 1.4-1.6 using the comparison formulas in Section 2.

Notation. All manifolds in this paper are smooth and without boundary. We denote by d the exterior differential operator and d^B when we like to insist the base manifold B.

We use the superconnection formalism of Quillen [32]. If A is a \mathbb{Z}_2 -graded algebra, and if $a, b \in A$, then we will note $[a, b] := ab - (-1)^{\deg(a) \deg(b)}ba$ as the supercommutator of a, b. If E, E' are two \mathbb{Z}_2 -graded spaces, we will note $E \widehat{\otimes} E'$ as the \mathbb{Z}_2 -graded tensor product as in [4, Section 1.3]. If one of E, E' is ungraded, we understand it as \mathbb{Z}_2 -graded by taking its odd part as zero.

For the fiber bundle $\pi : W \to B$, we use the sign convention for the integration of the differential forms along the oriented fibers Z as follows: for $\alpha \in \Omega^{\bullet}(B)$ and $\beta \in \Omega^{\bullet}(W),$

$$\int_{Z} (\pi^* \alpha) \wedge \beta = \alpha \wedge \int_{Z} \beta.$$
(1.15)

2 Bismut-Cheeger Eta Forms

In this section, we recall the definitions of the equivariant Bismut–Cheeger eta form and the equivariant version of the eta form with perturbation, which was originally introduced in [30, 31], and study the relations between them.

Let G be a compact Lie group. Let $\pi : W \to B$ be an equivariant submersion of two compact G-manifolds with compact fiber X such that G acts on B trivially.

Let $TX := \ker(\pi_* : TW \to TB)$ be the relative tangent bundle. Then TW and TX are equivariant vector bundles over W. Let $T^HW \subset TW$ be an equivariant horizontal subbundle, such that

$$TW = T^H W \oplus TX. \tag{2.1}$$

Since G is compact, such $T^H W$ always exists. Let $P^{TX} : TW \to TX$ be the projection associated with (2.1). Note that $T^H W \cong \pi^* TB$.

Let g^{TX} , g^{TB} be *G*-invariant metrics on *TX*, *TB*. We equip $TW = T^H W \oplus TX$ with the *G*-invariant metric $g^{TW} = \pi^* g^{TB} \oplus g^{TX}$. Let ∇^{TW} be the Levi–Civita connection on (TW, g^{TW}) . Let ∇^{TX} be the connection on *TX* defined by

$$\nabla^{TX} = P^{TX} \nabla^{TW} P^{TX}. \tag{2.2}$$

It is a *G*-invariant Euclidean connection on *TX* that depends only on $(T^H W, g^{TX})$ (cf. [6, Theorem 1.9]). Let ∇^{TB} be the Levi-Civita connection on (TB, g^{TB}) . Let $\nabla^{TB,TX}$ be the connection on *TW* defined by

$$\nabla^{TB,TX} = \pi^* \nabla^{TB} \oplus \nabla^{TX}, \tag{2.3}$$

which is also *G*-invariant.

Let C(TX) be the Clifford algebra bundle of (TX, g^{TX}) , whose fiber at $w \in W$ is the Clifford algebra $C(T_wX)$ of the Euclidean space (T_wX, g^{T_wX}) . We assume that TXhas a *G*-equivariant spin structure. Let S_X be the spinor bundle over W for (TX, g^{TX}) , which has a smooth action of C(TX) and is *G*-equivariant. Let ∇^{S_X} be the *G*-invariant connection on S_X induced by ∇^{TX} . If dim X is even, S_X is naturally \mathbb{Z}_2 -graded and ∇^{S_X} preserves this \mathbb{Z}_2 -grading.

Let E be an equivariant complex vector bundle over W. Let h^E be a G-invariant Hermitian metric on E. Let ∇^E be a G-invariant Hermitian connection on (E, h^E) . As in Section 1.1, we say $\underline{E} = (E, h^E, \nabla^E)$ is an equivariant geometric triple over W. Set $\nabla^{S_X \otimes E} := \nabla^{S_X} \otimes 1 + 1 \otimes \nabla^E$. Then $\nabla^{S_X \otimes E}$ is a G-invariant Hermitian connection on $(S_X \otimes E, h^{S_X} \otimes h^E)$.

Let $\{e_i\}_{i=1}^{\dim X}$ be a local orthonormal frame of TX. The fiberwise Dirac operator is defined by

$$D_X^E = \sum_{i=1}^{\dim X} c(e_i) \nabla_{e_i}^{\mathcal{S}_X \otimes E}, \qquad (2.4)$$

which is independent of the choice of the local orthonormal frame. If dim X is even, the fiberwise Dirac operator $D_X^E = D_+ \oplus D_-$ with respect to the \mathbb{Z}_2 -grading.

For $b \in B$, let \mathcal{E}_b be the set of smooth sections over $X_b = \pi^{-1}(b)$ of $\mathcal{S}_X \otimes E|_{X_b}$. As in [6], we will regard \mathcal{E} as an infinite dimensional vector bundle over B. If dim X is even, then \mathcal{E} is \mathbb{Z}_2 -graded.

Let $\nabla^{\mathcal{E}}$ be the connection on \mathcal{E} defined in [9, (1.7)], which preserves the L^2 inner product on \mathcal{E} . If $U \in TB$, let $U^H \in T^H W$ be its horizontal lift in $T^H W$ so that $\pi_* U^H = U$. Let $\{f_p\}$ be a local orthonormal frame of TB. We denote by $c(T) = -\frac{1}{2} c \left(P^{TX}[f_p^H, f_q^H] \right) f^p \wedge f^q \wedge$. Let \mathbb{B}_u be the rescaled Bismut superconnection defined by (see, e.g., [4, p. 336])

$$\mathbb{B}_u = \sqrt{u}D_X^E + \nabla^{\mathcal{E}} - \frac{c(T)}{4\sqrt{u}}, \quad u > 0.$$
(2.5)

Note that \mathbb{B}^2_u is a 2nd-order elliptic differential operator along the fibers [6, Theorem 3.5].

We assume that $\ker D_X^E$ forms a vector bundle over B. Then the L^2 inner product on \mathcal{E} induces a G-invariant metric on $\ker D_X^E$. Let $P^{\ker} : \mathcal{E} \to \ker D_X^E$ be the orthogonal projection with respect to the L^2 inner product on \mathcal{E} . Let

$$\nabla^{\ker} := P^{\ker} \circ \nabla^{\mathcal{E}} \circ P^{\ker}. \tag{2.6}$$

It is a *G*-invariant Hermitian connection on $\ker D_X^E$.

We define the supertrace Tr_s on a trace class operator on $\operatorname{End}(\mathcal{E})$ as in [4, Section 1.3]. If P is a trace class operator acting on $\Lambda(T^*B)\widehat{\otimes}\operatorname{End}(\mathcal{E})$ that takes values

in $\Lambda(T^*B)$, we use the convention that if $\omega \in \Lambda(T^*B)$,

$$\operatorname{Tr}[\omega P] = \omega \operatorname{Tr}[P], \quad \operatorname{Tr}_{s}[\omega P] = \omega \operatorname{Tr}_{s}[P].$$
 (2.7)

We denote by $\operatorname{Tr}^{\operatorname{odd/even}}[P]$ the part of $\operatorname{Tr}[P]$ that takes values in odd or even forms. Set

$$\widetilde{\mathrm{Tr}}[P] = \begin{cases} \mathrm{Tr}_{s}[P], & \text{if } \dim X \text{ is } \text{even}; \\ \mathrm{Tr}^{\mathrm{odd}}[P], & \text{if } \dim X \text{ is } \text{odd}. \end{cases}$$
(2.8)

For $\alpha \in \Omega^{j}(B)$, set

$$\psi_{B}(\alpha) = \begin{cases} \left(\frac{1}{2\pi\sqrt{-1}}\right)^{\frac{j}{2}} \cdot \alpha, & \text{if } j \text{ is even;} \\ \frac{1}{\sqrt{\pi}} \left(\frac{1}{2\pi\sqrt{-1}}\right)^{\frac{j-1}{2}} \cdot \alpha, & \text{if } j \text{ is odd.} \end{cases}$$
(2.9)

For $\beta \in \Omega^*(B \times [0, 1]_u)$, if $\beta = \beta_0 + du \wedge \beta_1$, with $\beta_0, \beta_1 \in \Lambda(T^*B)$, set

$$[\beta]^{du} := \beta_1. \tag{2.10}$$

The following definition is the equivariant version of the Bismut–Cheeger eta form in [8].

Definition 2.1. [21, Definition 2.3] For $g \in G$, the equivariant Bismut–Cheeger eta form is defined by

$$\begin{split} \tilde{\eta}_{g}^{BC}(\pi, E) &:= -\int_{0}^{\infty} \left\{ \psi_{\mathbb{R} \times B} \ \widetilde{\mathrm{Tr}} \left[g \exp\left(- \left(\mathbb{B}_{u} + \mathrm{d}u \wedge \frac{\partial}{\partial u} \right)^{2} \right) \right] \right\}^{du} \mathrm{d}u \\ &= \begin{cases} \int_{0}^{\infty} \frac{1}{\sqrt{\pi}} \psi_{B} \mathrm{Tr}^{\mathrm{even}} \left[g \ \frac{\partial \mathbb{B}_{u}}{\partial u} \exp(-\mathbb{B}_{u}^{2}) \right] \mathrm{d}u \in \Omega^{\mathrm{even}}(B, \mathbb{C}), \\ & \text{if } \dim X \text{ is } \mathrm{odd}; \\ \int_{0}^{\infty} \frac{1}{2\sqrt{\pi}\sqrt{-1}} \psi_{B} \mathrm{Tr}_{s} \left[g \ \frac{\partial \mathbb{B}_{u}}{\partial u} \exp(-\mathbb{B}_{u}^{2}) \right] \mathrm{d}u \in \Omega^{\mathrm{odd}}(B, \mathbb{C}), \\ & \text{if } \dim X \text{ is } \mathrm{even}. \end{cases} \end{split}$$
(2.11)

Note that the convergence of the integrals in (11) are highly nontrivial [21, (2.72), (2.77)] (cf. also [4, Theorem 10.32]). Remark that by changing the variable (see also

[21, Remark 2.4]),

$$\tilde{\eta}_{g}^{BC}(\pi, E) = -\int_{0}^{\infty} \left\{ \psi_{\mathbb{R} \times B} \, \widetilde{\mathrm{Tr}} \left[g \exp\left(-\left(\mathbb{B}_{u^{2}} + du \wedge \frac{\partial}{\partial u} \right)^{2} \right) \right] \right\}^{du} \, \mathrm{d}u.$$
(2.12)

We define the reduced equivariant eta form $\tilde{\tilde{\eta}}_{g}^{BC}(\pi, E)$ as in (1.9). If $E = E_{+} \oplus E_{-}$ is \mathbb{Z}_{2} -graded, we define $\tilde{\tilde{\eta}}_{g}^{BC}(\pi, E) = \tilde{\tilde{\eta}}_{g}^{BC}(\pi, E_{+}) - \tilde{\tilde{\eta}}_{g}^{BC}(\pi, E_{-})$.

Now, we recall the definition of the equivariant eta form with perturbation in [23].

Let $K_G^0(B)$ be the equivariant topological K^0 -group of B. Fix $s \in S^1$. Let $\iota : B \to B \times S^1$, $b \mapsto (b,s)$, be the *G*-equivariant inclusion map, where we suppose that the *G*-action on S^1 is trivial. From [33, Definitions 2.7 and 2.8],

$$K^1_G(B) \simeq \ker \left(\iota^* : K^0_G(B \times S^1) \to K^0_G(B)\right).$$
(2.13)

Recall that *G* acts on *B* trivially. For $x \in K_G^0(B)$, $g \in G$, the classical equivariant Chern character map sends *x* to $ch_g(x) \in H^{even}(B, \mathbb{C})$. By (2.13), for $x \in K_G^1(B)$, we can regard *x* as an element x' in $K_G^0(B \times S^1)$. The odd equivariant Chern character map

$$ch_{g}: K^{1}_{G}(B) \longrightarrow H^{odd}(B, \mathbb{C})$$
(2.14)

is defined by $\operatorname{ch}_g(x) := \int_{S^1} \operatorname{ch}_g(x')$.

If dim X is even (resp. odd), the equivariant (analytic) index $\operatorname{Ind}(D_X^E) \in K_G^0(B)$ (resp. $K_G^1(B)$). If $\operatorname{Ind}(D_X^E) = 0 \in K_G^*(B)$, by [23, Proposition 3.3 (i)] (cf. also [30, Proposition 1] and [31, Proposition 2]), there exists a smooth family of equivariant bounded pseudodifferential operators \mathcal{A} such that $(D_X^E + \mathcal{A})|_{X_b}$ is invertible for any $b \in B$. If dim X is even, \mathcal{A} is additionally required to anti-commute with the \mathbb{Z}_2 -grading of the spinor \mathcal{S}_X . (If $E = E_+ \oplus E_-$ is \mathbb{Z}_2 -graded, \mathcal{A} is required to anti-commute with the \mathbb{Z}_2 -grading of $\mathcal{S}_X \widehat{\otimes} E$ when dim X is even and commute with the \mathbb{Z}_2 -grading of $\mathcal{S}_X \widehat{\otimes} E$ when dim X is called a perturbation operator.

Let $\chi\in \mathcal{C}^\infty(\mathbb{R})$ be a cut-off function such that

$$\chi(u) = \begin{cases} 0, & \text{if } u < 1; \\ 1, & \text{if } u > 2. \end{cases}$$
(2.15)

Set

$$\mathbb{B}'_{u} = \mathbb{B}_{u} + \sqrt{u}\chi(\sqrt{u})\mathcal{A}.$$
(2.16)

The following definition is the equivariant version of the eta form with perturbation in [17, 30, 31].

Definition 2.2. [23, Definition 3.12] Modulo exact forms on *B*, the equivariant eta form with perturbation operator A is defined by

$$\begin{split} \tilde{\eta}_{g}(\pi,\mathcal{A}) &:= -\int_{0}^{\infty} \left\{ \psi_{\mathbb{R}\times B} \, \widetilde{\mathrm{Tr}} \left[g \exp\left(- \left(\mathbb{B}'_{u} + du \wedge \frac{\partial}{\partial u} \right)^{2} \right) \right] \right\}^{du} \mathrm{d}u \\ &= \left\{ \begin{array}{l} \int_{0}^{\infty} \, \frac{1}{\sqrt{\pi}} \psi_{B} \mathrm{Tr}^{\mathrm{even}} \left[g \, \frac{\partial \mathbb{B}'_{u}}{\partial u} \exp(-(\mathbb{B}'_{u})^{2}) \right] \mathrm{d}u, \quad \mathrm{if} \, \mathrm{dim} \, X \, \mathrm{is} \, \mathrm{odd}; \\ \int_{0}^{\infty} \, \frac{1}{2\sqrt{\pi}\sqrt{-1}} \psi_{B} \mathrm{Tr}_{s} \left[g \, \frac{\partial \mathbb{B}'_{u}}{\partial u} \exp(-(\mathbb{B}'_{u})^{2}) \right] \mathrm{d}u, \quad \mathrm{if} \, \mathrm{dim} \, X \, \mathrm{is} \, \mathrm{even}. \end{split} \right.$$

Moreover, as in (1.8) (see, e.g., [23, (3.66)]),

$$d\tilde{\eta}_g(\pi, \mathcal{A}) = \int_{X^g} \widehat{A}_g(TX, \nabla^{TX}) \mathrm{ch}_g(E, \nabla^E).$$
(2.18)

As in (2.12), we have

$$\tilde{\eta}_{g}(\pi, \mathcal{A}) = -\int_{0}^{\infty} \left\{ \psi_{\mathbb{R} \times B} \, \widetilde{\mathrm{Tr}} \left[g \exp\left(- \left(\mathbb{B}'_{u^{2}} + du \wedge \frac{\partial}{\partial u} \right)^{2} \right) \right] \right\}^{du} \, \mathrm{d}u.$$
(2.19)

The proof of the well-definedness of $\tilde{\eta}_g(\pi, \mathcal{A})$ is the same as that of $\tilde{\eta}_g^{BC}(\pi, E)$. Remark that if we choose another cut-off function, the difference of the new eta form and the original one is an exact form. So modulo exact forms, the definition of $\tilde{\eta}_g(\pi, \mathcal{A})$ is independent of the choice of the cut-off function in (2.15).

Now, we discuss the relations between these two types of equivariant eta forms when ker D_X^E forms a vector bundle over *B*.

Case 1: dim X is odd.

Since $\ker D_X^E$ forms a vector bundle, $\ker(D_X^E + P^{\ker}) = 0$. So from [23, Proposition 3.3] (cf. also [30, Proposition 1]), we have $\operatorname{Ind}(D_X^E) = 0 \in K_G^1(B)$ and P^{\ker} is a perturbation operator.

Proposition 2.3. If dim X is odd, we have

$$\tilde{\eta}_g(\pi, P^{\text{ker}}) = \tilde{\tilde{\eta}}_g^{BC}(\pi, E) \mod d\Omega^*(B).$$
(2.20)

Proof. Let $\widetilde{\mathbb{B}}_u$ be the Bismut superconnection with respect to the fiber bundle $W \times [0,1]_s \to B \times [0,1]_s$ with fiber X such that restricted on $W \times \{s\}, s \in [0,1], \widetilde{\mathbb{B}}_u|_{W \times \{s\}} = \mathbb{B}_u + ds \wedge \frac{\partial}{\partial s}$. Set

$$\widetilde{\mathbb{B}}'_{u^2} := \widetilde{\mathbb{B}}_{u^2} + us\chi(u)P^{\mathrm{ker}} = u(D_X^E + s\chi(u)P^{\mathrm{ker}}) + \nabla^{\mathcal{E}} + ds \wedge \frac{\partial}{\partial s} - \frac{c(T)}{4u}.$$
(2.21)

We decompose

$$\psi_{\mathbb{R}^{2}\times B} \operatorname{Tr}^{\operatorname{odd}} \left[g \exp\left(-\left(\widetilde{\mathbb{B}}_{u^{2}}^{\prime} + du \wedge \frac{\partial}{\partial u} \right)^{2} \right) \right]$$
$$= du \wedge \gamma(u, s) + ds \wedge r_{1}(u, s) + du \wedge ds \wedge r_{2}(u, s) + r_{3}(u, s), \quad (2.22)$$

where γ , r_1 , r_2 , r_3 do not contain du neither ds. From (2.12), (2.19), and (22), we have

$$\tilde{\eta}_g^{BC}(\pi, E) = -\int_0^\infty \gamma(u, 0) \,\mathrm{d}u, \quad \tilde{\eta}_g(\pi, P^{\mathrm{ker}}) = -\int_0^\infty \gamma(u, 1) \,\mathrm{d}u. \tag{2.23}$$

By [4, Theorem 9.17],

$$\left(du \wedge \frac{\partial}{\partial u} + ds \wedge \frac{\partial}{\partial s} + d^{B}\right) \operatorname{Tr}\left[g \exp\left(-\left(\widetilde{\mathbb{B}}'_{u^{2}} + du \wedge \frac{\partial}{\partial u}\right)^{2}\right)\right] = 0.$$
(2.24)

So from (2.22) and (2.24),

$$\frac{\partial \gamma(u,s)}{\partial s} = \frac{\partial r_1(u,s)}{\partial u} + d^B r_2(u,s).$$
(2.25)

From (2.23) and (2.25),

$$\tilde{\eta}_{g}(\pi, P^{\text{ker}}) - \tilde{\eta}_{g}^{BC}(\pi, E) = -\int_{0}^{+\infty} (\gamma(u, 1) - \gamma(u, 0)) \, du$$

$$= -\int_{0}^{+\infty} \int_{0}^{1} \frac{\partial}{\partial s} \gamma(u, s) \, ds \, du = -\int_{0}^{1} \int_{0}^{+\infty} \frac{\partial}{\partial s} \gamma(u, s) \, du \, ds$$

$$= -\int_{0}^{1} \int_{0}^{+\infty} \frac{\partial}{\partial u} r_{1}(u, s) \, du \, ds - d^{B} \int_{0}^{1} \int_{0}^{+\infty} r_{2}(u, s) \, du \, ds$$

$$= \int_{0}^{1} \left(r_{1}(0, s) - r_{1}(\infty, s) \right) \, ds - d^{B} \int_{0}^{1} \int_{0}^{+\infty} r_{2}(u, s) \, du \, ds. \quad (2.26)$$

The commutative property of the integrals in the above formula is guaranteed by [21, (2.72), (2.77)] for $s \in [0, 1]$. By (2.15), (2.22), and the equivariant family local index theorem (see, e.g., [25, Theorem 1], [21, Theorem 2.2]), as in [21, (2.96), (2.97)], we have

$$r_1(0,s) = 0, \quad r_1(\infty,s) = \frac{1}{\sqrt{\pi}} \lim_{u \to +\infty} \psi_B \left\{ \operatorname{Tr}^{\mathrm{odd}} \left[g \exp\left(- \left(\widetilde{\mathbb{B}}'_{u^2} \right)^2 \right) \right] \right\}^{ds}.$$
(2.27)

Therefore, by (2.26) and (2.27), modulo exact forms on *B*, we have

$$\tilde{\eta}_{g}(\pi, P^{\text{ker}}) - \tilde{\eta}_{g}^{BC}(\pi, E) = -\frac{1}{\sqrt{\pi}} \lim_{u \to +\infty} \psi_{B} \int_{0}^{1} \left\{ \text{Tr}^{\text{odd}} \left[g \exp\left(- \left(\widetilde{\mathbb{B}}'_{u^{2}} \right)^{2} \right) \right] \right\}^{ds} \, \mathrm{d}s.$$
(2.28)

For a family of bounded operators \mathcal{A}_u , $u \in \mathbb{R}_+$, we write $\mathcal{A}_u = O(u^{-k})$ as $u \to +\infty$ if there exists C > 0 such that if u is large enough, the norm of \mathcal{A}_u is dominated by Cu^{-k} .

Let $P^{\ker,\perp}$ be the orthogonal projection on the orthogonal complement of $\ker D^E_X$. Set

$$\begin{split} E_{u} &= P^{\mathrm{ker}} \circ \left(\widetilde{\mathbb{B}}'_{u^{2}}\right)^{2} \circ P^{\mathrm{ker}}, \qquad F_{u} = P^{\mathrm{ker}} \circ \left(\widetilde{\mathbb{B}}'_{u^{2}}\right)^{2} \circ P^{\mathrm{ker},\perp}, \\ G_{u} &= P^{\mathrm{ker},\perp} \circ \left(\widetilde{\mathbb{B}}'_{u^{2}}\right)^{2} \circ P^{\mathrm{ker}}, \qquad H_{u} = P^{\mathrm{ker},\perp} \circ \left(\widetilde{\mathbb{B}}'_{u^{2}}\right)^{2} \circ P^{\mathrm{ker},\perp}. \end{split}$$
(2.29)

Note that

$$P^{\mathrm{ker}}\{u(D_X^E + s\chi(u)P^{\mathrm{ker}})\}P^{\mathrm{ker}} = us\chi(u)P^{\mathrm{ker}}.$$
(2.30)

Denote by $\widetilde{
abla} =
abla^{\mathcal{E}} + ds \wedge rac{\partial}{\partial s}$. Let

$$E' = u^{2}s^{2}P^{\text{ker}} + uds \wedge P^{\text{ker}} + P^{\text{ker}}(\nabla^{\mathcal{E}})^{2}P^{\text{ker}},$$

$$F' = P^{\text{ker},\perp}[D_{X}^{E} + sP^{\text{ker}}, \widetilde{\nabla}]P^{\text{ker}},$$

$$G' = P^{\text{ker}}[D_{X}^{E} + sP^{\text{ker}}, \widetilde{\nabla}]P^{\text{ker},\perp},$$

$$H' = P^{\text{ker},\perp}(D_{X}^{E} + sP^{\text{ker}})^{2}P^{\text{ker},\perp} = P^{\text{ker},\perp}(D_{X}^{E})^{2}P^{\text{ker},\perp}.$$
(2.31)

By (2.21), when $u
ightarrow +\infty$,

$$\begin{split} E_u &= E' + O(u^{-1}), \qquad F_u = uF' + F'' + O(1), \\ G_u &= uG' + G'' + O(1), \quad H_u = u^2H' + uH'' + H''' + O(1), \end{split} \tag{2.32}$$

where F'', G'', H'', H''' are 1st-order differential operators along the fiber. Moreover, by (2.6) and (2.31), we have

$$E' - G'H'^{-1}F' = u^2s^2P^{\mathrm{ker}} + uds \wedge + (\nabla^{\mathrm{ker}})^2 + s^2\mathcal{C} + s\mathcal{D}, \qquad (2.33)$$

where

$$\mathcal{C} = P^{\ker \widetilde{\nabla}} P^{\ker,\perp} (P^{\ker,\perp} D_X^E P^{\ker,\perp})^{-2} P^{\ker,\perp} \widetilde{\nabla} P^{\ker},$$

$$\mathcal{D} = 2P^{\ker \widetilde{\nabla}} P^{\ker,\perp} (P^{\ker,\perp} D_X^E P^{\ker,\perp})^{-1} P^{\ker,\perp} \widetilde{\nabla} P^{\ker}.$$
(2.34)

Following the same way as the proof of [21, Theorem 5.13] (cf. also [4, Theorem 9.19]), when $u \to +\infty$, we can obtain

$$\exp\left(-\left(\widetilde{\mathbb{B}}'_{u^{2}}\right)^{2}\right) = P^{\ker} \circ \exp\left(-\left(E' - G'H'^{-1}F'\right)\right) \circ P^{\ker} + O(u^{-1}).$$
(2.35)

So from (2.28), as in [21, Theorem 5.15], modulo exact forms,

$$\tilde{\eta}_{g}(\pi, P^{\text{ker}}) - \tilde{\eta}_{g}^{BC}(\pi, E) = \frac{1}{\sqrt{\pi}} \lim_{u \to +\infty} \int_{0}^{1} u e^{-u^{2}s^{2}} \psi_{B} \text{Tr}^{\text{even}}[g \exp(-(\nabla^{\text{ker}})^{2} - s^{2}\mathcal{C} - s\mathcal{D})] \, \mathrm{d}s$$
$$= \frac{1}{\sqrt{\pi}} \lim_{u \to +\infty} \int_{0}^{u} e^{-t^{2}} \psi_{B} \text{Tr}^{\text{even}}[g \exp(-(\nabla^{\text{ker}})^{2} - u^{-2}t^{2}\mathcal{C} - u^{-1}t\mathcal{D})] \, \mathrm{d}t. \quad (2.36)$$

Using the Volterra series (cf., e.g., [4, (2.5)]), we have

$$\tilde{\eta}_g(\pi, P^{\mathrm{ker}}) - \tilde{\eta}_g^{BC}(\pi, E) = \frac{1}{\sqrt{\pi}} \int_0^{+\infty} e^{-t^2} \,\mathrm{d}t \cdot \mathrm{ch}_g(\mathrm{ker}\, D_X^E, \nabla^{\mathrm{ker}}) = \frac{1}{2} \mathrm{ch}_g(\mathrm{ker}\, D_X^E, \nabla^{\mathrm{ker}}). \tag{2.37}$$

By (1.9) and (2.37), the proof of Proposition 2.3 is completed.

Remark that we can also obtain this proposition using the method in [4, Chapter 9] as in [30, Proposition 17].

Remark 2.4. (a) Observe that from (1.8) and (2.18),

$$d^{B}\tilde{\eta}_{g}(\pi, P^{\text{ker}}) = d^{B}\tilde{\tilde{\eta}}_{g}^{BC}(\pi, E) = \int_{X^{g}} \widehat{A}_{g}(TX, \nabla^{TX}) \text{ch}_{g}(E, \nabla^{E}), \qquad (2.38)$$

which is compatible with this proposition.

(b) If B is a point, from [23, Remark 3.20], both sides of (2.20) are equal to the reduced eta invariant.

Case 2: dim X is even.

In this case, $\ker D_X^E = \ker D_+ \oplus \ker D_-$ is a \mathbb{Z}_2 -graded vector bundle. But P^{\ker} is not a perturbation operator because P^{\ker} does not anti-commute with the \mathbb{Z}_2 -grading of the spinor. In general, $\operatorname{Ind}(D_X^E) \neq 0 \in K_G^0(B)$.

We consider a new equivariant fiber bundle $\pi' : W \sqcup B \to B$ with fiber $X \sqcup \{\text{pt}\}$, which is the disjoint union of the fiber bundle $W \to B$ and a new fiber bundle $B \times \{\text{pt}\} \to B$ whose fiber is a point. We denote by $(\ker D_+)_-$ and $(\ker D_-)_+$ the vector bundles $\ker D_+$ and $\ker D_-$ over $B \times \{\text{pt}\}$ with the inverse grading. Let D^H be the Dirac operator on $H = \mathcal{E} \oplus (\ker D_+)_- \oplus (\ker D_-)_+$ with respect to the fiber bundle π' . Then $D^H = D_X^E \oplus 0 \oplus 0$. Note that

$$\operatorname{Ind}(D^{H}) = \operatorname{Ind}D_{X}^{E} - \operatorname{Ind}D_{X}^{E} = 0 \in K_{G}^{0}(B).$$
(2.39)

Let $\tilde{\eta}_g^{BC}(\pi', E \sqcup (\ker D_X^E)^{op})$ be the equivariant Bismut–Cheeger eta form associated with this fiber bundle π' . Here, $(\ker D_X^E)^{op}$ denotes the \mathbb{Z}_2 -graded vector bundle $\ker D_X^E$ with the opposite grading. Notice that if the fiber is a point, the corresponding Bismut– Cheeger eta form vanishes. Thus, we have

$$\tilde{\eta}_g^{BC}(\pi', E \sqcup (\ker D_X^E)^{op}) = \tilde{\eta}_g^{BC}(\pi, E).$$
(2.40)

Let $V = \ker D_+ \oplus \ker D_- \oplus (\ker D_+)_- \oplus (\ker D_-)_+$. Let $P^V : H \to V$ be the orthogonal projection on V. Let

$$\mathcal{A} = \left(\begin{array}{cc} 0 & 1\\ 1 & 0 \end{array}\right) \circ P^V \tag{2.41}$$

on $(\ker D_+ \oplus (\ker D_-)_+) \oplus ((\ker D_+)_- \oplus \ker D_-)$. Then \mathcal{A} anticommutes with the \mathbb{Z}_2 -grading and $\mathcal{A}^2 = \text{Id on } V$. Since

$$(D^{H} + A)^{2} = D^{H,2} + A^{2} = D^{H,2} + P^{V},$$
(2.42)

we see that $D^H + A$ is invertible. So A is a perturbit operator associated with D^H . Let $\tilde{\eta}_g^H(\pi', A)$ be the corresponding eta form with perturbation operator A.

Proposition 2.5. If dim X is even, we have

$$\tilde{\eta}_g^H(\pi', \mathcal{A}) = \tilde{\tilde{\eta}}_g^{BC}(\pi, E) \mod d\Omega^*(B).$$
(2.43)

Proof. From (1.9) and (2.40), we only need to prove $\tilde{\eta}_g^H(\pi', \mathcal{A}) = \tilde{\eta}_g^{BC}(\pi', E \sqcup (\ker D_X^E)^{op}) \mod d\Omega^*(B).$

Let $\nabla^{H} = \nabla^{\mathcal{E}} \oplus \nabla^{\text{ker}}$ on *H*. As in (2.21), we set

$$\widetilde{\mathbb{B}}_{u^{2}}^{\prime H} = u(D^{H} + s\chi(u)\mathcal{A}) + \nabla^{H} + ds \wedge \frac{\partial}{\partial s} - \frac{c(T)}{4u}$$
(2.44)

with respect to the vector bundle $(W \sqcup B) \times [0, 1] \rightarrow B \times [0, 1]$ with fiber $X \sqcup \{pt\}$.

Then following the same process as in (2.22)– (2.28), modulo exact forms on B, we have

$$\begin{split} \tilde{\eta}_{g}(\pi',\mathcal{A}) &- \tilde{\eta}_{g}^{BC}(\pi',E \sqcup (\ker D_{X}^{E})^{op}) \\ &= -\frac{1}{2\sqrt{\pi}\sqrt{-1}} \lim_{u \to +\infty} \psi_{B} \int_{0}^{1} \left\{ \operatorname{Tr}_{s} \left[g \exp\left(-\left(\widetilde{\mathbb{B}}_{u^{2}}^{\prime H}\right)^{2}\right) \right] \right\}^{ds} \mathrm{d}s. \quad (2.45) \end{split}$$

Let $R^{\text{ker}} = (\nabla^{\text{ker}})^2$ be the curvature of the connection ∇^{ker} on $\text{ker} D_X^E$. As in (2.29)– (2.37), modulo exact forms on B,

$$\begin{split} \tilde{\eta}_{g}(\pi',\mathcal{A}) &- \tilde{\eta}_{g}^{BC}(\pi',E \sqcup (\ker D_{X}^{E})^{op}) \\ &= \frac{1}{2\sqrt{\pi}\sqrt{-1}}\psi_{B}\int_{0}^{+\infty}e^{-t^{2}}\,\mathrm{d}t\cdot\mathrm{Tr}_{s}^{V}[g\mathcal{A}\exp(-R^{\ker}\oplus R^{\ker})]. \end{split}$$
(2.46)

Note that

$$\mathcal{A} \exp(-R^{\ker} \oplus R^{\ker}) = \begin{pmatrix} 0 & 1\\ 1 & 0 \end{pmatrix} \begin{pmatrix} \exp(-R^{\ker}) & 0\\ 0 & \exp(-R^{\ker}) \end{pmatrix}$$
$$= \begin{pmatrix} 0 & \exp(-R^{\ker})\\ \exp(-R^{\ker}) & 0 \end{pmatrix}. \quad (2.47)$$

Since the group action commutes with \mathcal{A} and R^{ker} , $\operatorname{Tr}_{s}^{V}[g\mathcal{A}\exp(-R^{\text{ker}}\oplus R^{\text{ker}})] = 0$. So Proposition 2.5 follows from (2.40), (2.46), and (2.47).

The proof of Proposition 2.5 is completed.

Remark 2.6. As in (2.38), from (1.8) and (2.18), we have

$$d^{B}\tilde{\eta}_{g}^{H}(\pi',\mathcal{A}) = d^{B}\tilde{\tilde{\eta}}_{g}^{BC}(\pi',E) = \int_{X^{g}} \widehat{A}_{g}(TX,\nabla^{TX}) \operatorname{ch}_{g}(E,\nabla^{E}) - \operatorname{ch}_{g}(\ker D_{X}^{E},\nabla^{\ker}), \quad (2.48)$$

which is compatible with Proposition 2.5.

3 Anomaly Formula

In this section, we will prove Theorem 1.4. We use the notations and the assumptions in Section 1.3.

Let $s \in I$, I = [0, 1], parametrize a smooth path of equivariant horizontal subbundles $\{T_s^H W\}_{s \in [0,1]}$. The existence of such path follows from the affineness of the space of equivariant splitting maps in the short exact sequence $0 \rightarrow TX \rightarrow TW \rightarrow \pi^*TB \rightarrow 0$. Let g_s^{TX} and h_s^E be *G*-invariant metrics on *TX* and *E* connecting (g_0^{TX}, h_0^E) and (g_1^{TX}, h_1^E) smoothly. Let ∇_s^E be a *G*-invariant Hermitian connection on (E, h_s^E) connecting ∇_0^E and ∇_1^E smoothly. Now, we get a smooth path connecting the geometric data $(T_0^H W, g_0^{TX}, h_0^E, \nabla_0^E)$ and $(T_1^H W, g_1^{TX}, h_1^E, \nabla_1^E)$. Let $D_{X,s}^E$ be the fiberwise Dirac operator associated with $(T_s^H W, g_s^{TX}, \nabla_s^E)$.

If $\operatorname{Ind}(D_{X,0}^E) = 0 \in K_G^*(B)$, by the homotopy invariance of the equivariant family index, we have $\operatorname{Ind}(D_{X,1}^E) = 0 \in K_G^*(B)$. In this case, let \mathcal{A}_0 , \mathcal{A}_1 be perturbation operators with respect to $D_{X,0}^E$, $D_{X,1}^E$, respectively. Let $\tilde{\eta}_{g,0}(\pi, \mathcal{A}_0)$ and $\tilde{\eta}_{g,1}(\pi, \mathcal{A}_1)$ be corresponding equivariant eta forms with perturbations. Let P_0 , P_1 be orthogonal projections onto the eigenspaces of the positive spectrum of $D_{X,0}^E + \mathcal{A}_0$ and $D_{X,1}^E + \mathcal{A}_1$, respectively. Let $\mathrm{sf}_G\{(D_{X,0}^E + \mathcal{A}_0, P_0), (D_{X,1}^E + \mathcal{A}_1, P_1)\} \in K_G^*(B)$ be the equivariant Dai–Zhang higher spectral flow defined in [23, Definitions 3.7 and 3.8] (cf. also [17, Definition 1.5]), which we denote by $\mathrm{sf}_G\{D_{X,0}^E + \mathcal{A}_0, D_{X,1}^E + \mathcal{A}_1\}$ for simplicity. The following proposition was established in [23], which is a generalization of [17, Theorem 0.1].

Proposition 3.1. [23, Theorem 1.2] For any $g \in G$, modulo exact forms on *B*, we have

$$\begin{split} \tilde{\eta}_{g,1}(\pi,\mathcal{A}_1) &- \tilde{\eta}_{g,0}(\pi,\mathcal{A}_0) = \int_{X^g} \widetilde{\widehat{A}}_g(TX,\nabla_0^{TX},\nabla_1^{TX}) \operatorname{ch}_g(E,\nabla_0^E) \\ &+ \int_{X^g} \widehat{A}_g(TX,\nabla_1^{TX}) \widetilde{\operatorname{ch}}_g(E,\nabla_0^E,\nabla_1^E) + \operatorname{ch}_g\left(\operatorname{sf}_G\{D_{X,0}^E + \mathcal{A}_0, D_{X,1}^E + \mathcal{A}_1\}\right). \end{split}$$
(3.1)

Case 1: dim *X* is odd.

In this case, Theorem 1.4 follows directly from Propositions 2.3 and 3.1.

Theorem 3.2. If dim X is odd, for any $g \in G$, modulo exact forms on *B*, we have

$$\begin{split} \tilde{\eta}_{g,1}^{BC}(\pi,E) &- \tilde{\eta}_{g,0}^{BC}(\pi,E) = \int_{X^g} \widetilde{\widehat{A}}_g(TX,\nabla_0^{TX},\nabla_1^{TX}) \operatorname{ch}_g(E,\nabla_0^E) \\ &+ \int_{X^g} \widehat{A}_g(TX,\nabla_1^{TX}) \widetilde{\operatorname{ch}}_g(E,\nabla_0^E,\nabla_1^E) + \operatorname{ch}_g\left(\operatorname{sf}_G\{D_{X,0}^E + P_0^{\operatorname{ker}}, D_{X,1}^E + P_1^{\operatorname{ker}}\}\right), \end{split}$$
(3.2)

Here, P_0^{ker} and P_1^{ker} are the orthogonal projections onto $\ker D_{X,0}^E$ and $\ker D_{X,1}^E$, respectively.

Remark that if there is a smooth path connecting $(T_0^H W, g_0^{TX}, h_0^E, \nabla_0^E)$ and $(T_1^H W, g_1^{TX}, h_1^E, \nabla_1^E)$ such that dim ker $D_{X,s}^E$ is locally constant, we have sf_G $\{D_{X,0}^E + P_0^{\text{ker}}, D_{X,1}^E + P_1^{\text{ker}}\} = 0$ and $[ch(\ker D_{X,0}^E)] = [ch(\ker D_{X,1}^E)]$. This is the case in [21, Theorem 1.2].

Case 2: dim X is even.

In this case, recall that we assume $\ker D_{X,i}^E = \ker D_{+,i} \oplus \ker D_{-,i}$, i = 0, 1, are \mathbb{Z}_2 -graded vector bundles over *B*. Remark that if $s \in (0, 1)$, dim $\ker D_{X,s}^E$ is usually not locally constant and $\ker D_{X,0}^E$ is not isomorphic to $\ker D_{X,1}^E$ as equivariant vector bundles. We will use a perturbation trick in [19, 26] to overcome this difficulty.

Consider the equivariant fiber bundle $\pi \times \text{Id} : W \times [0,1] \to B \times [0,1]$ and the equivariant \mathbb{Z}_2 -graded Hilbert bundle \mathcal{H}_{\pm} over $B \times [0,1]$ with fiber $L^2(X_{(b,s)}, \mathcal{S}_{\pm}(TX) \otimes E|_{X_{(b,s)}})$ for $(b,s) \in B \times [0,1]$. Let $T^H \widetilde{W}|_{B \times \{s\}} := T_s^H W \oplus \mathbb{R}$, which is an equivariant horizontal subbundle with respect to $\pi \times \text{Id}$. Let $(g^{TX}, h^{\tilde{E}}, \nabla^{\tilde{E}})$ be the geometric data associated with $\pi \times \text{Id}$ induced from $(g_s^{TX}, h_s^E, \nabla_s^E)$. Let $\tilde{D}_X^E = \tilde{D}_+ \oplus \tilde{D}_-$ be the corresponding \mathbb{Z}_2 -graded fiberwise Dirac operator. Then $\tilde{D}_{\pm}|_{B \times \{0\}} = D_{\pm,0}$ and $\tilde{D}_{\pm}|_{B \times \{1\}} = D_{\pm,1}$.

By [19, Lemma 7.13] (see also [26, p.27]), there exist finite dimensional vector subbundles $L_{\pm} \subset \mathcal{H}_{\pm}$ and complementary closed subbundles $K_{\pm} \subset \mathcal{H}_{\pm}$ over $B \times [0, 1]$, that is, $\mathcal{H}_{\pm} = L_{\pm} \oplus K_{\pm}$, such that $\tilde{D}_{\pm} \in \operatorname{Hom}(\mathcal{H}_{\pm}, \mathcal{H}_{\mp})$ is block diagonal as a map $\tilde{D}_{\pm} : L_{\pm} \oplus K_{\pm} \to L_{\mp} \oplus K_{\mp}$ and \tilde{D}_{\pm} restricts to an isomorphism between K_{+} and K_{-} . Remark that the direct sums here may not be orthogonal. It is easy to see that we could choose L_{\pm} and K_{\pm} to be *G*-equivariant.

Note that ker $\tilde{D}_{\pm} \subset L_{\pm}$. Add a new \mathbb{Z}_2 -graded vector bundle $(L_+)_- \oplus (L_-)_+$ over $B \times [0,1] \times \{\text{pt}\}$. As in Section 2, $(L_+)_-$ and $(L_-)_+$ are vector bundles L_+ and L_- with the inverse grading.

As in Section 2, we consider the equivariant fiber bundle $\pi' \times \mathrm{Id} : (W \sqcup B) \times [0, 1] \to B \times [0, 1]$ with fiber $X \sqcup \{\mathrm{pt}\}$. Let \tilde{D}^L be the fiberwise Dirac operator on $\mathcal{H}^L = \mathcal{H} \oplus (L_+)_- \oplus (L_-)_+$. Then $\tilde{D}^L = \tilde{D}^E_X \oplus 0 \oplus 0$. Let $P^L : \mathcal{H}^L \to L_+ \oplus L_- \oplus (L_+)_- \oplus (L_-)_+$ be the projection. Let

$$\mathcal{A}^{L} = \left(\begin{array}{cc} 0 & 1\\ 1 & 0 \end{array}\right) \circ P^{L} \tag{3.3}$$

on $(L_+ \oplus (L_-)_+) \oplus ((L_+)_- \oplus L_-)$. From [19, Lemma 7.20], $\tilde{D}^L + \mathcal{A}^L$ is invertible. Since \mathcal{A}^L anti-commutes with the \mathbb{Z}_2 -grading of \mathcal{H}^L , it is a perturbation operator associated with \tilde{D}^L .

For i = 0, 1, ker $D_{\pm,i}$ are equivariant subbundles of $L_{\pm,i} := L_{\pm}|_{B \times \{i\}}$. Let $F_{\pm,i}$ be the orthogonal complements of ker $D_{\pm,i}$ in $L_{\pm,i}$. Then $L_{\pm,i} = \ker D_{\pm,i} \oplus F_{\pm,i}$. Since $D_{\pm,i} : F_{\pm,i} \to F_{\pm,i}$, we have $F_{\pm,i} \simeq F_{-,i}$ as equivariant vector bundles over $B \times \{i\}$, i = 0, 1. We identify $F_{\pm,i}$ and $F_{-,i}$ via the isomorphism $D_{\pm,i}$. Let $P_i^F : \mathcal{H}^L|_{B \times \{i\}} \to \ker D_{X,i}^E \oplus (L_{\pm,i})_{\pm} \oplus (L_{\pm,i})_{\pm}$ be the orthogonal projection. Let

$$\mathcal{A}_{F,i}^{\mathrm{ker}} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \circ P_i^F$$
(3.4)

on

$$\left(\ker D_{+,i} \oplus (\ker D_{-,i})_+\right) \oplus \left((\ker D_{+,i})_- \oplus \ker D_{-,i}\right) \oplus (F_{-,i})_+ \oplus (F_{+,i})_-.$$

Let $D_i^L = \tilde{D}^L|_{B \times \{i\}}$, i = 0, 1. Then $\mathcal{A}_{F,i}^{\text{ker}}$ are perturbation operators associated with D_i^L .

For i = 0, 1, let $\nabla_i^{\mathcal{E}}$ and ∇_i^{ker} be connections in (2.5) and (2.6) associated with $(T_i^H W, g_i^{TX}, h_i^E, \nabla_i^E)$. Let ∇_i^L and ∇_i^F be the projected compatible connections on L_i and F_i , respectively, as in (2.6). Let $\nabla_i^{\mathcal{E},F} = \nabla_i^{\mathcal{E}} \oplus \nabla_i^{\text{ker}} \oplus \nabla_i^F$ on $\mathcal{H}^L|_{B \times \{i\}}$. Set

$$\widetilde{\mathbb{B}}_{u^{2},i}^{F} := u(D_{i}^{L} + \chi(u)\mathcal{A}_{F,i}^{\ker}) + \nabla_{i}^{\mathcal{E},F} - \frac{c(T_{i})}{4u},$$
(3.5)

where $c(T_i)$ is associated with $(T_i^H W, g_i^{TX})$. Let $\tilde{\eta}_{g,i}^H(\pi', \mathcal{A}_{F,i}^{\text{ker}})$ be corresponding eta forms for $g \in G$.

For two equivariant geometric triples $\underline{E} = (E, h^E, \nabla^E)$ and $\underline{F} = (F, h^F, \nabla^F)$, if $E \simeq F$ as equivariant vector bundles over B, we denote by $\widetilde{ch}_g(\underline{E}, \underline{F})$ the equivariant Chern–Simons form between \underline{E} and \underline{F} for $g \in G$. Moreover, we have

$$d \widetilde{\mathrm{ch}}_g(\underline{E}, \underline{F}) = \mathrm{ch}_g(F, \nabla^F) - \mathrm{ch}_g(E, \nabla^E). \tag{3.6}$$

Lemma 3.3. If dim *X* is even, for $g \in G$, i = 0, 1, we have

$$\tilde{\eta}_{g,i}^{H}(\pi', \mathcal{A}_{F,i}^{\mathrm{ker}}) = \tilde{\tilde{\eta}}_{g,i}^{BC}(\pi, E) - \widetilde{\mathrm{ch}}_{g}\left(\underline{F_{-,i}}, \underline{F_{+,i}}\right) \mod d\Omega^{*}(B).$$
(3.7)

Proof. For i = 0, 1, let

$$\mathcal{A}_i = \left(\begin{array}{cc} 0 & 1\\ 1 & 0 \end{array}\right) \tag{3.8}$$

on $(\ker D_{+,i} \oplus (\ker D_{-,i})_{+}) \oplus ((\ker D_{+,i})_{-} \oplus \ker D_{-,i})$. Let

$$\mathcal{A}_{F,i} = \left(\begin{array}{cc} 0 & 1\\ 1 & 0 \end{array}\right) \tag{3.9}$$

on $(F_{-,i})_+ \oplus (F_{+,i})_-$. Then $\mathcal{A}_{F,i}^{\mathrm{ker}} = (\mathcal{A}_i \oplus \mathcal{A}_{F,i}) \circ P_i^F$. By Proposition 2.5, we have

$$\tilde{\eta}_g^H(\pi', \mathcal{A}_i) = \tilde{\tilde{\eta}}_{g,i}^{BC}(\pi, E) \mod d\Omega^*(B).$$
(3.10)

From Definition 2.2, (2.12), and (2.19), modulo exact forms on *B*,

$$\tilde{\eta}_{g,i}^{H}(\pi',\mathcal{A}_{F,i}^{\mathrm{ker}}) = \tilde{\eta}_{g}^{H}(\pi',\mathcal{A}_{i}) - \int_{0}^{\infty} \frac{\sqrt{-1}}{2\pi} \mathrm{Tr}_{s} \left[g \, \frac{\partial \nabla_{u^{2},i}^{F}}{\partial u} \exp\left(-\frac{(\nabla_{u^{2},i}^{F})^{2}}{2\pi\sqrt{-1}}\right) \right] \mathrm{d}u, \qquad (3.11)$$

where $\nabla_{u,i}^F = \nabla_{+,i}^F \oplus \nabla_{-,i}^F + \sqrt{u}\chi(\sqrt{u})\mathcal{A}_{F,i}$ on $F_{+,i} \oplus F_{-,i}$. We denote by

$$\hat{\eta}_{g}(\nabla_{u^{2},i}^{F}) := \int_{0}^{\infty} \frac{\sqrt{-1}}{2\pi} \operatorname{Tr}_{s} \left[g \, \frac{\partial \nabla_{u^{2},i}^{F}}{\partial u} \, \exp\!\left(-\frac{(\nabla_{u^{2},i}^{F})^{2}}{2\pi\sqrt{-1}}\right) \right] \mathrm{d}u, \tag{3.12}$$

which is the equivariant eta form with perturbation when the fiber is a point and is the analogue of the analytic torsion for a finite dimensional complex [10]. Let $\widetilde{\nabla}_{u,i}^F := \nabla_{-,i}^F \oplus \nabla_{-,i}^F + \sqrt{u}\mathcal{A}_{F,i}$ on $(F_{-,i})_+ \oplus F_{-,i}$. We deform the metric and the connection from $F_{+,i} \oplus F_{-,i}$ to $F_{-,i} \oplus F_{-,i}$. Since $\mathcal{A}_{F,i}$ is invertible, by [7, Theorem 2.10], modulo exact forms, we have

$$\hat{\eta}_{g}(\nabla_{u^{2},i}^{F}) - \hat{\eta}_{g}(\widetilde{\nabla}_{u^{2},i}^{F}) = \widetilde{\operatorname{ch}}_{g}\left(\underbrace{(F_{-,i})_{+}}_{+} \oplus \underbrace{F_{-,i}}_{+,i} \oplus \underbrace{F_{-,i}}_{+,i} \oplus \underbrace{F_{-,i}}_{-,i}\right) = \widetilde{\operatorname{ch}}_{g}\left(\underbrace{F_{-,i}}_{+,i}, \underbrace{F_{+,i}}_{+,i}\right).$$
(3.13)

From [7, (2.24)], we see that $\hat{\eta}_g(\widetilde{\nabla}^F_{u^2,i}) = 0$. So modulo exact forms, by (3.13), we have

$$\hat{\eta}_g(\nabla^F_{u^2,i}) = \widetilde{\mathrm{ch}}_g\left(\underline{F_{-,i}}, \underline{F_{+,i}}\right).$$
(3.14)

From (3.10)– (3.12) and (3.14), we obtain Lemma 3.1.

Since *L* is an equivariant vector bundle over $B \times [0, 1]$, L_0 is isomorphic to L_1 as equivariant \mathbb{Z}_2 -graded vector bundles over *B*.

Theorem 3.4. If dim X is even, for any $g \in G$, modulo exact forms on *B*, we have

$$\begin{split} \tilde{\eta}_{g,1}^{BC}(\pi, E) &- \tilde{\eta}_{g,0}^{BC}(\pi, E) = \int_{X^g} \widetilde{\mathbf{A}}_g(TX, \nabla_0^{TX}, \nabla_1^{TX}) \operatorname{ch}_g(E, \nabla_0^E) \\ &+ \int_{X^g} \widehat{\mathbf{A}}_g(TX, \nabla_1^{TX}) \widetilde{\mathbf{Ch}}_g(E, \nabla_0^E, \nabla_1^E) + \operatorname{ch}_g \left(\operatorname{sf}_G \{ D_0^L + \mathcal{A}_{F,0}^{\operatorname{ker}}, D_1^L + \mathcal{A}_{F,1}^{\operatorname{ker}} \} \right) \\ &+ \widetilde{\mathbf{Ch}}_g(\underline{F}_{-,1}, \underline{F}_{+,1}) - \widetilde{\mathbf{Ch}}_g(\underline{F}_{-,0}, \underline{F}_{+,0}) - \widetilde{\mathbf{Ch}}_g(\underline{L}_0, \underline{L}_1) \\ &+ \widetilde{\mathbf{Ch}}_g(L_1, \nabla_1^{\operatorname{ker}} \oplus \nabla_1^F, \nabla_1^L) - \widetilde{\mathbf{Ch}}_g(L_0, \nabla_0^{\operatorname{ker}} \oplus \nabla_0^F, \nabla_0^L). \end{split}$$
(3.15)

Proof. For i = 0, 1, let $\nabla_i^{\mathcal{E},L} = \nabla_i^{\mathcal{E}} \oplus \nabla_i^{L}$. Denote by $\mathcal{A}_i^{L} = \mathcal{A}^{L}|_{B \times \{i\}}$. Set

$$\widetilde{\mathbb{B}}_{u^{2},i}^{L} := u(D_{i}^{L} + \chi(u)\mathcal{A}_{i}^{L}) + \nabla_{i}^{\mathcal{E},L} - \frac{c(T_{i})}{4u}.$$
(3.16)

Let $\tilde{\eta}_{g,i}^H(\pi', \mathcal{A}_i^L)$ be corresponding equivariant eta forms. Since $\mathrm{sf}_G\{D_0^L + \mathcal{A}_0^L, D_1^L + \mathcal{A}_1^L\} = 0$, from Proposition 3.1, modulo exact forms, we have

$$\begin{split} \tilde{\eta}_{g,1}^{H}(\pi',\mathcal{A}_{1}^{L}) &- \tilde{\eta}_{g,0}^{H}(\pi',\mathcal{A}_{0}^{L}) = \int_{X^{g}} \widehat{A}_{g}(TX,\nabla_{0}^{TX},\nabla_{1}^{TX}) \operatorname{ch}_{g}(E,\nabla_{0}^{E}) \\ &+ \int_{X^{g}} \widetilde{\widehat{A}}_{g}(TX,\nabla_{1}^{TX}) \widetilde{\operatorname{ch}}_{g}(E,\nabla_{0}^{E},\nabla_{1}^{E}) - \widetilde{\operatorname{ch}}_{g}(\underline{L}_{0},\underline{L}_{1}) \quad (3.17) \end{split}$$

and

$$\begin{split} \tilde{\eta}_{g,i}^{H}(\pi',\mathcal{A}_{i}^{L}) &- \tilde{\eta}_{g,i}^{H}(\pi',\mathcal{A}_{F,i}^{\ker}) = -\widetilde{\mathrm{ch}}_{g}(L_{i},\nabla_{i}^{\ker} \oplus \nabla_{i}^{F},\nabla_{i}^{L}) \\ &+ \mathrm{ch}_{g}(\mathrm{sf}_{G}\{D_{i}^{L} + \mathcal{A}_{F,i}^{\ker},D_{i}^{L} + \mathcal{A}_{i}^{L}\}). \end{split}$$
(3.18)

Since

$$sf_{G}\{D^{L} + \mathcal{A}_{F,0}^{ker}, D^{L} + \mathcal{A}_{F,1}^{ker}\} = sf_{G}\{D^{L} + \mathcal{A}_{F,0}^{ker}, D^{L} + \mathcal{A}_{0}^{L}\} - sf_{G}\{D^{L} + \mathcal{A}_{F,1}^{ker}, D^{L} + \mathcal{A}_{1}^{L}\}, \quad (3.19)$$

Theorem 3.4 follows from (3.17), (3.18), and Lemma 3.1.

The proof of Theorem 3.4 is completed.

Remark that by [24, Definition 2.14], for any equivariant Chern–Simons form $\widetilde{ch}_g(\cdot)$, $[0, \widetilde{ch}_g(\cdot)] \in \widehat{K}_g^0(B)$. So, by [24, Proposition 3.1], $\widetilde{ch}_g(\cdot) \in ch_g(K_G^1(B))$. (In [24], it is easy to see that Proposition 3.1 holds for any compact Lie group *G* when *G* acts on *B* trivially.) Therefore, we obtain Theorem 1.4.

The following lemma will be used in Sections 4 and 5.

Lemma 3.5. Assume that dim X is even and $\operatorname{Ind}(D_X^E) = 0 \in K_G^0(B)$. Let \mathcal{A}_X be a perturbation operator associated with D_X^E . Then there exists $x \in K_G^1(B)$, which can be constructed as an equivariant higher spectral flow explicitly, such that modulo exact forms, for $g \in G$,

$$\widetilde{\eta}_g(\pi, \mathcal{A}_X) = \widetilde{\widetilde{\eta}}_g^{BC}(\pi, E) + \mathrm{ch}_g(x) \mod d\Omega^*(B).$$
(3.20)

Proof. Since $\operatorname{Ind}(D_X^E) = 0 \in K_G^0(B)$, by [33, Proposition 2.4], there exists $n \in \mathbb{N}$, such that $\ker D_+ \oplus \varepsilon_+^n \simeq \ker D_- \oplus \varepsilon_-^n$, where ε_{\pm}^n are equivariant trivial *n*-dimensional complex vector bundles over *B* with \mathbb{Z}_2 -grading. We identify $\ker D_+ \oplus \varepsilon_+^n$ and $\ker D_- \oplus \varepsilon_-^n$ via this isomorphism. We consider ε_{\pm}^n as vector bundles over $B \times \{\text{pt}\}$ with trivial connections.

Let $P_{\varepsilon}: \mathcal{E} \oplus \varepsilon_{+}^{n} \oplus \varepsilon_{-}^{n} \to \ker D_{X}^{E} \oplus \varepsilon_{+}^{n} \oplus \varepsilon_{-}^{n}$ be the projection. Let

$$\mathcal{A}_{\varepsilon}^{\mathrm{ker}} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \circ P_{\varepsilon} \tag{3.21}$$

on $(\ker D_+ \oplus \varepsilon_+^n) \oplus (\ker D_- \oplus \varepsilon_-^n)$. Let

$$\mathcal{A}_{\varepsilon} = \begin{pmatrix} \mathcal{A}_{X} & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \circ P_{\varepsilon}$$
(3.22)

on $\mathcal{E} \oplus \varepsilon_+^n \oplus \varepsilon_-^n$. Then $\mathcal{A}_{\varepsilon}^{\text{ker}}$ and $\mathcal{A}_{\varepsilon}$ are perturbation operators associated with the fiberwise Dirac operator $D_X^H = D_X^E \oplus 0 \oplus 0$ on $\mathcal{E} \oplus \varepsilon_+^n \oplus \varepsilon_-^n$. Let $\tilde{\eta}_g^H(\pi', \mathcal{A}_{\varepsilon}^{\text{ker}})$ and $\tilde{\eta}_g^H(\pi', \mathcal{A}_{\varepsilon})$ be corresponding equivariant eta forms with perturbation. By the proof of Proposition 2.5, we have

$$\widetilde{\eta}_{g}^{H}(\pi', \mathcal{A}_{\varepsilon}^{\ker}) = \widetilde{\widetilde{\eta}}_{g}^{BC}(\pi, E) \mod d\Omega^{*}(B).$$
(3.23)

From (3.11) and (3.14), since the connections on ε^n_\pm are trivial, modulo exact forms, we have

$$\widetilde{\eta}_{g}(\pi, \mathcal{A}_{X}) = \widetilde{\eta}_{g}^{H}(\pi', \mathcal{A}_{\varepsilon}).$$
(3.24)

From Proposition 3.1, modulo exact forms, we have

$$\widetilde{\eta}_{g}^{H}(\pi',\mathcal{A}_{\varepsilon}) - \widetilde{\eta}_{g}^{H}(\pi',\mathcal{A}_{\varepsilon}^{\mathrm{ker}}) = \mathrm{ch}_{g}(\mathrm{sf}\{D_{X}^{H} + \mathcal{A}_{\varepsilon}^{\mathrm{ker}}, D_{X}^{H} + \mathcal{A}_{\varepsilon}\}).$$
(3.25)

So from (3.23)–(3.25), modulo exact forms, we have

$$\widetilde{\eta}_{g}(\pi, \mathcal{A}_{X}) - \mathrm{ch}_{g}(\mathrm{sf}\{D_{X}^{H} + \mathcal{A}_{\varepsilon}^{\mathrm{ker}}, D_{X}^{H} + \mathcal{A}_{\varepsilon}\}) = \widetilde{\tilde{\eta}}_{g}^{BC}(\pi, E) \mod d\Omega^{*}(B).$$
(3.26)

The proof of Lemma 3.5 is completed.

4 Embedding Formula

In this section, we will prove Theorem 1.5. We use the notations and the assumptions in Section 1.4.

Let $\xi = \xi_+ \oplus \xi_-$ be the \mathbb{Z}_2 -graded vector bundle over W. From [22, Proposition 3.6], if $\operatorname{Ind} D_Y^{\mu} = 0 \in K_G^*(B)$, then $\operatorname{Ind} D_X^{\xi} = 0 \in K_G^*(B)$. Moreover, we have the following proposition, which is the equivariant family extension of the main result in [11].

Proposition 4.1. [22, Theorem 3.7] Assume that $\operatorname{Ind} D_Y^{\mu} = 0 \in K_G^*(B)$. Let \mathcal{A}_Y and \mathcal{A}_X be perturbation operators associated with D_Y^{μ} and D_X^{ξ} , respectively. Then there exists a perturbation operator $\mathcal{A}_{\mathcal{A}_Y,X}$ associated with D_X^{ξ} , depending on \mathcal{A}_Y , such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\eta}_{g}(\pi_{X},\mathcal{A}_{X}) &= \tilde{\eta}_{g}(\pi_{Y},\mathcal{A}_{Y}) + \int_{X^{g}} \widehat{A}_{g}(TX,\nabla^{TX}) \gamma_{g}^{X}(\mathcal{F}_{Y},\mathcal{F}_{X}) \\ &+ \int_{Y^{g}} \widetilde{A}_{g}(\nabla^{TY,N},\nabla^{TX|_{W^{g}}}) \, \widehat{A}_{g}^{-1}(N,\nabla^{N}) \operatorname{ch}_{g}(\mu,\nabla^{\mu}) \\ &+ \operatorname{ch}_{g}(\operatorname{sf}_{G}\{D_{X}^{\xi} + \mathcal{A}_{\mathcal{A}_{Y},X}, D_{X}^{\xi} + \mathcal{A}_{X}\}). \end{split}$$
(4.1)

Remark that in [22], the author assumed that the embedding is totally geodesic. In fact, this condition is not necessary because we can use the variation formula Proposition 3.1 to obtain the formula for the general embedding from the totally geodesic case.

Case 1: dim X is odd.

In this case, dim Y is also odd. Let P_Y^{ker} and P_X^{ker} be orthogonal projections onto $\ker D_Y^{\mu}$ and $\ker D_X^{\xi}$, respectively. So, from Proposition 2.3, we have

$$\begin{split} \tilde{\eta}_g(\pi_Y, P_Y^{\text{ker}}) &= \tilde{\tilde{\eta}}_g^{BC}(\pi_Y, \mu) \mod d\Omega^*(B), \\ \tilde{\eta}_g(\pi_X, P_X^{\text{ker}}) &= \tilde{\tilde{\eta}}_g^{BC}(\pi_X, \xi) \mod d\Omega^*(B). \end{split}$$
(4.2)

Theorem 4.2. If dim X is odd, modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{X},\xi) &= \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{Y},\mu) + \int_{X^{g}} \widehat{A}_{g}(TX,\nabla^{TX}) \gamma_{g}^{X}(\mathcal{F}_{Y},\mathcal{F}_{X}) \\ &+ \int_{Y^{g}} \widetilde{\tilde{A}}_{g}(\nabla^{TY,N},\nabla^{TX|_{W^{g}}}) \widehat{A}_{g}^{-1}(N,\nabla^{N}) \operatorname{ch}_{g}(\mu,\nabla^{\mu}) \\ &+ \operatorname{ch}_{g}(\operatorname{sf}_{G}\{D_{X}^{\xi} + \mathcal{A}_{P_{v}^{\ker},X}, D_{X}^{\xi} + P_{X}^{\ker}\}). \end{split}$$
(4.3)

Case 2: dim X is even.

In this case, dim Y is even. We denote by $D_Y^{\mu} = D_{Y,+} \oplus D_{Y,-}$. As in the setting of Proposition 2.5, we add $B \times \{\text{pt}\}$ and consider the embedding $W \sqcup B \to V \sqcup B$. We denote by $(\ker D_{Y,-})_+$ and $(\ker D_{Y,+})_-$ the vector bundles $\ker D_{Y,-}$ and $\ker D_{Y,+}$ over $B \times \{\text{pt}\}$ with the inverse grading and assume that restricted on $B \times \{\text{pt}\}$, the embedding is the identity map. This new embedding also forms an equivariant version of the Atiyah–Hirzebruch direct image from $\mu \sqcup (\ker D_Y^{\mu})^{op}$ to $\xi \sqcup (\ker D_Y^{\mu})^{op}$.

Let $\pi'_W : W \sqcup B \to B$ be the fiber bundle induced from π_W . Let P^Y be the projection onto $\ker D^{\mu}_V \oplus (\ker D^{\mu}_Y)^{op}$. Set

$$\mathcal{A}_Y = \left(\begin{array}{cc} 0 & 1\\ 1 & 0 \end{array}\right) \circ P^Y \tag{4.4}$$

on $(\ker D_{Y,+} \oplus (\ker D_{Y,-})_+) \oplus ((\ker D_{Y,+})_- \oplus \ker D_{Y,-})$. Then, from Proposition 2.5, modulo exact forms on B,

$$\widetilde{\eta}_{g}^{H}(\pi'_{W},\mathcal{A}_{Y}) = \widetilde{\widetilde{\eta}}_{g}^{BC}(\pi_{W},\mu).$$
(4.5)

Let $\pi'_V : V \sqcup B \to B$ be the fiber bundle induced from π_V . Let $\tilde{\eta}_g^{BC}(\pi'_V, \xi \sqcup (\ker D_Y^{\mu})^{op})$ be the equivariant Bismut–Cheeger eta form associated with this fiber bundle. Since the equivariant Bismut–Cheeger eta form vanishes when the fiber is a point, as in (2.40), we have

$$\tilde{\eta}_g^{BC}(\pi_V,\xi) = \tilde{\eta}_g^{BC}(\pi'_V,\xi \sqcup (\ker D_Y^{\mu})^{op}).$$

$$(4.6)$$

Let $D^{H,\mu}$ and $D^{H,\xi}$ be fiberwise Dirac operators associated with π'_W and π'_V . Since $\operatorname{Ind}(D^{H,\mu}) = 0 \in K^0_G(B)$, $\operatorname{Ind}(D^{H,\xi}) = 0 \in K^0_G(B)$. Let \mathcal{A}_X be a perturbation operator associated with $D^{H,\xi}$. From Lemma 3.5, there exists $x \in K^1_G(B)$, such that modulo exact forms,

$$\tilde{\eta}_{g}^{BC}(\pi'_{V},\xi \sqcup (\ker D_{Y}^{\mu})^{op}) = \tilde{\eta}_{g}(\pi'_{V},\mathcal{A}_{X}) - \mathrm{ch}_{g}(x).$$

$$(4.7)$$

From (1.9), (4.5)–(4.7), and Proposition 4.1, we obtain Theorem 1.5 when dim X is even.

Theorem 4.3. If dim X is even, there exists $x \in K^1_G(B)$ such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\eta}_{g}^{BC}(\pi_{V},\xi_{+}) &- \tilde{\eta}_{g}^{BC}(\pi_{V},\xi_{-}) = \tilde{\eta}_{g}^{BC}(\pi_{W},\mu) + \int_{X^{g}} \widehat{A}_{g}(TX,\nabla^{TX})\gamma_{g}(Y^{g},X^{g}) \\ &+ \int_{Y^{g}} \widetilde{\widehat{A}}_{g}(TY,\nabla^{TY,N},\nabla^{TX}|_{W^{g}}) \widehat{A}_{g}^{-1}(N,\nabla^{N}) \mathrm{ch}_{g}(E,\nabla^{E}) + \mathrm{ch}_{g}(x). \end{split}$$
(4.8)

5 Functoriality

In this section, we will prove Theorem 1.6. We use the notations and the assumptions in Section 1.5.

From [23, Lemma 3.6], if $\operatorname{Ind} D_X^E = 0 \in K_G^*(V)$, then $\operatorname{Ind} D_Z^E = 0 \in K_G^*(B)$. In this case, the functoriality of the equivariant eta forms with perturbation was obtained in [23].

Proposition 5.1. [23, Theorem 1.3] Let \mathcal{A}_Z and \mathcal{A}_X be perturbation operators associated with D_Z^E and D_X^E . Then there exists $x \in K_G^*(B)$, an equivariant higher spectral flow, such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \widetilde{\eta}_{g}(\pi_{3},\mathcal{A}_{Z}) &= \int_{Y^{g}} \widehat{A}_{g}(TY,\nabla^{TY}) \, \widetilde{\eta}_{g}(\pi_{1},\mathcal{A}_{X}) \\ &+ \int_{Z^{g}} \widetilde{\widehat{A}}_{g}(TZ,\nabla^{TY,TX},\nabla^{TZ}) \operatorname{ch}_{g}(E,\nabla^{E}) + \operatorname{ch}_{g}(x). \end{split}$$
(5.1)

Case 1: dim *X* is odd and dim *Y* is even. Then dim *Z* is odd.

Since ker D_X^E and ker D_Z^E form equivariant vector bundles, $\operatorname{Ind} D_X^E = \operatorname{Ind} D_Z^E = 0 \in K_G^1(B)$. Let $P^{\ker D_X}$ and $P^{\ker D_Z}$ be orthogonal projections on ker D_X^E and ker D_Z^E . Then $P^{\ker D_X}$

and $P^{\ker D_Z}$ are perturbation operators associated with D_X^E and D_Z^E , respectively, as in Section 2. By Proposition 2.3,

$$\begin{split} \tilde{\eta}_g(\pi_1, P^{\ker D_X}) &= \tilde{\tilde{\eta}}_g^{BC}(\pi_1, E) \mod d\Omega^*(V), \\ \tilde{\eta}_g(\pi_3, P^{\ker D_Z}) &= \tilde{\tilde{\eta}}_g^{BC}(\pi_3, E) \mod d\Omega^*(B). \end{split}$$
(5.2)

From Proposition 5.1 and (5.2), we obtain the functoriality in this case.

Theorem 5.2. If dim X is odd and dim Y is even, there exists $x \in K^0_G(B)$, such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\eta}_{g}^{BC}(\pi_{3}, E) &= \int_{Y^{g}} \widehat{A}_{g}(TY, \nabla^{TY}) \, \tilde{\eta}_{g}^{BC}(\pi_{1}, E) \\ &+ \int_{Z^{g}} \widetilde{\widehat{A}}_{g}(TZ, \nabla^{TY, TX}, \nabla^{TZ}) \operatorname{ch}_{g}(E, \nabla^{E}) + \operatorname{ch}_{g}(x). \end{split}$$
(5.3)

Case 2: dim X is odd and dim Y is odd. Then dim Z is even.

In this case, $\operatorname{Ind} D_X^E = 0 \in K_G^1(V)$ and $P^{\ker D_X}$ is a perturbation operator. From [23, Lemma 3.6], $\operatorname{Ind} D_Z^E = 0 \in K_G^0(B)$. Let \mathcal{A}_Z be a perturbation operator with respect to D_Z^E . Then, by Lemma 3.5, there exists $x \in K_G^1(B)$, such that modulo exact forms on B, for $g \in G$,

$$\widetilde{\eta}_g(\pi_3, \mathcal{A}_Z) = \widetilde{\widetilde{\eta}}_g^{BC}(\pi_3, E) + \mathrm{ch}_g(x) \mod d\Omega^*(B).$$
(5.4)

By Proposition 2.3,

$$\tilde{\eta}_g(\pi_1, P^{\ker D_X}) = \tilde{\tilde{\eta}}^{BC}(\pi_1, E) \mod d\Omega^*(V).$$
(5.5)

So from Proposition 5.1, (5.4), and (5.5), we have the following theorem.

Theorem 5.3. If dim X and dim Y are both odd, there exists $x \in K^1_G(B)$, such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\eta}_{g}^{BC}(\pi_{3}, E) &= \int_{Y^{g}} \widehat{A}_{g}(TY, \nabla^{TY}) \, \tilde{\eta}_{g}^{BC}(\pi_{1}, E) \\ &+ \int_{Z^{g}} \widetilde{\widehat{A}}_{g}(TZ, \nabla^{TY, TX}, \nabla^{TZ}) \operatorname{ch}_{g}(E, \nabla^{E}) + \operatorname{ch}_{g}(x). \end{split}$$
(5.6)

Case 3: dim X is even and dim Y is odd. Then dim Z is odd.

Recall that $\ker D_X^E = \ker D_{X,+} \oplus \ker D_{X,-}$ is a \mathbb{Z}_2 -graded equivariant vector bundle over V. We consider a new equivariant fiber bundle $\pi'_1 : W \sqcup V \to V$ with fiber $X \sqcup \{\text{pt}\}$ and with vector bundle $(\ker D_{X,-})_+ \oplus (\ker D_{X,+})_-$ over $V \times \{\text{pt}\}$. Let $\tilde{\eta}_g^H(\pi'_1, \mathcal{A})$ be the equivariant eta form with perturbation \mathcal{A} in (2.41) on $(\ker D_{X,+} \oplus (\ker D_{X,-})_+) \oplus ((\ker D_{X,+})_- \oplus \ker D_{X,-})$. By Proposition 2.5, we have

$$\tilde{\eta}_g^H(\pi_1', \mathcal{A}) = \tilde{\tilde{\eta}}_g^{BC}(\pi_1, E) \mod d\Omega^*(V).$$
(5.7)

We consider the equivariant fiber bundle $\pi_3 \sqcup \pi_2^- : W \sqcup V \to B$ such that the geometric triple over W is \underline{E} and that over V is $(\underline{\ker D_X^E})^{op}$. Since $\ker D_Y^{\underline{\ker}D_X^E}$ and $\ker D_Z^E$ are equivariant vector bundles, from Definition 2.1,

$$\widetilde{\eta}_g^{BC}(\pi_3 \sqcup \pi_2^-, E \sqcup (\ker D_X^E)^{op}) = \widetilde{\eta}_g^{BC}(\pi_3, E) - \widetilde{\eta}_g^{BC}(\pi_2, \ker D_X^E).$$
(5.8)

Since dim Z and dim Y are both odd, by Proposition 2.3, there exists a perturbation operator \mathcal{A}_P of the fiberwise Dirac operator associated with the fiber bundle $\pi_3 \sqcup \pi_2^-$, such that modulo exact forms, for $g \in G$,

$$\begin{split} \widetilde{\eta}_{g}(\pi_{3} \sqcup \pi_{2}^{-}, \mathcal{A}_{P}) &= \widetilde{\eta}_{g}^{BC}(\pi_{3} \sqcup \pi_{2}^{-}, E \sqcup (\ker D_{X}^{E})^{op}) \\ &+ \frac{1}{2} \mathrm{ch}_{g}(\ker D_{Z}^{E}) - \frac{1}{2} \mathrm{ch}_{g}(\ker D_{Y}^{E}). \end{split}$$
(5.9)

So from Proposition 5.1 and (5.7)–(5.9), we have the following theorem.

Theorem 5.4. If dim X is even and dim Y is odd, there exists $x \in K^0_G(B)$, such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{3}, E) &= \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{2}, \ker D_{X}^{E}) + \int_{Y^{g}} \widehat{A}_{g}(TY, \nabla^{TY}) \, \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{1}, E) \\ &+ \int_{Z^{g}} \widetilde{\tilde{A}}_{g}(TZ, \nabla^{TY, TX}, \nabla^{TZ}) \operatorname{ch}_{g}(E, \nabla^{E}) + \operatorname{ch}_{g}(x). \end{split}$$
(5.10)

Case 4: $\dim X$ is even and $\dim Y$ is even. Then $\dim Z$ is even.

We consider the fiber bundles in Case 3. In this case, (5.7) and (5.8) still hold. Let \tilde{D}_X be the fiberwise Dirac operator with respect to π'_1 . Then $\operatorname{Ind}(\tilde{D}_X) = 0 \in K^0_G(V)$. By [23, Lemma 3.6], $\operatorname{Ind}D_Z^E - \operatorname{Ind}D_Y^{\ker D_X^E} = 0 \in K^0_G(B)$. By Lemma 3.5, in this case, (9) is replaced by

$$\widetilde{\eta}_g(\pi_3 \sqcup \pi_2^-, \mathcal{A}_P) = \widetilde{\eta}_g^{BC}(\pi_3 \sqcup \pi_2^-, E \sqcup (\ker D_X^E)^{op}) + \mathrm{ch}_g(x).$$
(5.11)

Here, *x* is an element of $K_G^1(B)$.

Theorem 5.5. If dim X and dim Y are both even, there exists $x \in K^1_G(B)$, such that modulo exact forms on B, for $g \in G$,

$$\begin{split} \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{3},E) &= \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{2},\ker D_{X}^{E}) + \int_{Y^{g}} \widehat{A}_{g}(TY,\nabla^{TY}) \, \tilde{\tilde{\eta}}_{g}^{BC}(\pi_{1},E) \\ &+ \int_{Z^{g}} \widetilde{A}_{g}(TZ,\nabla^{TY,TX},\nabla^{TZ}) \operatorname{ch}_{g}(E,\nabla^{E}) + \operatorname{ch}_{g}(\mathbf{x}). \end{split}$$
(5.12)

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